

Some conclusions on fluid-structure interactions

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1. Well-posedness and input-output stability
2. Asymptotic behaviour

Well-posedness and input-output stability

The space $\mathcal{H}^\infty(\mathbb{C}_0, W)$ consist of all the analytic functions $G : \mathbb{C}_0 \rightarrow Z$ for which

$$\sup_{s \in \mathbb{C}_0} \|G(s)\|_W < \infty.$$

The condition $G \in \mathcal{H}^\infty(\mathbb{C}_0, L(U, Y))$ is equivalent to the fact that if $u \in L^2([0, \infty); U)$, then $y \in L^2([0, \infty); Y)$.

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We consider the linear model

$$\frac{\partial h}{\partial t} + \frac{\partial q}{\partial x} = 0, \quad (x \in \mathcal{E}), \quad (1)$$

$$\frac{\partial q}{\partial t} + \frac{\partial h}{\partial x} - \mu \frac{\partial^2 q}{\partial x^2} = 0, \quad (x \in \mathcal{E}), \quad (2)$$

$$\dot{h}_S(t) + \frac{\partial q}{\partial x} = 0 \quad (x \in \mathcal{I}), \quad (3)$$

$$\frac{\partial q}{\partial t} + \frac{\partial p}{\partial x} = 0 \quad (x \in \mathcal{I}), \quad (4)$$

$$h(t, a^-) - \mu \frac{\partial q}{\partial x}(t, a^-) = p(t, a^+) + h_S(t) - \mu \frac{\partial q}{\partial x}(t, a^+), \quad (5)$$

$$h(t, b^+) - \mu \frac{\partial q}{\partial x}(t, b^+) = p(t, b^-) + h_S(t) - \mu \frac{\partial q}{\partial x}(t, b^-), \quad (6)$$

$$\ddot{h}_S(t) = \int_a^b p(t, x) dx + u(t) \quad (t > 0). \quad (7)$$

Here, we consider a output given

$$y(t) = h_S(t) \quad (t \geq 0). \quad (8)$$

Our first main result is a following reformulation of the system.

Set

$$X := \mathbb{C} \times H^1(\mathcal{E}) \times L^2(\mathcal{E}) \times \mathbb{C} \times \mathbb{C}. \quad (9)$$

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Theorem

Equations (1)-(8) can be recast as

$$\begin{aligned} \dot{z} &= Az + Bu \\ y &= Cz, \end{aligned} \tag{10}$$

where the components of the vector $z(t)$ are $h_S(t)$, $h(t, \cdot)$, $q(t, \cdot)$, $q(t, a)$ and $q(t, b)$, B is in $\mathcal{L}(\mathbb{C}, X)$, C is in $\mathcal{L}(X, \mathbb{C})$ and A is a generator of an analytic semigroup on X .

Corollary

Equations (1)-(8) define a well-posed linear system with state space X defined in (9) and input and output spaces $U = Y = \mathbb{C}$.

Informally, this means: on any time interval $[T, t]$, for any initial state x_0 and any input function u , it has a unique state trajectory x and a unique output function y , both defined on $[T, t]$.

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For $t \geq 0$, we set $q_a(t) := q(t, a)$ and $q_b(t) := q(t, b)$. Since (3) implies that q is a linear function of x on \mathcal{I} , for every $t \geq 0$ and $x \in \mathcal{I}$,

$$\dot{h}_S(t) = -\frac{q_b(t) - q_a(t)}{b - a}, \quad (11)$$

$$q(t, x) = q_a(t) \left(\frac{x - b}{a - b} \right) + q_b(t) \left(\frac{x - a}{b - a} \right), \quad (12)$$

$$\frac{\partial q}{\partial x}(t, x) = \frac{q_b(t) - q_a(t)}{b - a}. \quad (13)$$

We differentiate (4) with respect to x and use (5)-(3) to arrive at



$$\begin{aligned} \frac{\partial^2 p}{\partial x^2}(t, x) &= \ddot{h}_S(t) \quad (x \in \mathcal{I}), \\ p(t, a^+) &= p_a(t), \quad p(t, b^-) = p_b(t), \end{aligned} \quad (14)$$

where

$$p_a(t) := h(t, a^-) - \mu \frac{\partial q}{\partial x}(t, a^-) - h_{sol}(t) - \mu \dot{h}_{sol}(t), \quad (15)$$

$$p_b(t) := h(t, b^+) - \mu \frac{\partial q}{\partial x}(t, b^+) - h_{sol}(t) - \mu \dot{h}_{sol}(t). \quad (16)$$

The first equation in (14) implies that, for every $t \geq 0$, $p(t, x)$ is a second order polynomial in x so that

$$\begin{aligned}\int_a^b p(t, x) dx &= p(t, a)l - \dot{q}_a(t)\frac{l^2}{3} - \dot{q}_b(t)\frac{l^2}{6} \\ &= p(t, b)l + \dot{q}_a(t)\frac{l^2}{6} + \dot{q}_b(t)\frac{l^2}{3},\end{aligned}$$

where we set $l := b - a$.

Combining this with (7) and (11) we deduce that



$$\begin{aligned} \left[1 + \frac{l^3}{3}\right] \dot{q}_a(t) - \left[1 - \frac{l^3}{6}\right] \dot{q}_b(t) &= p(t, a)l^2 + lu(t), \\ - \left[1 - \frac{l^3}{6}\right] \dot{q}_a(t) + \left[1 + \frac{l^3}{3}\right] \dot{q}_b(t) &= -p(t, b)l^2 - lu(t). \end{aligned}$$

Inverting the above linear system, we get

$$\begin{bmatrix} \dot{q}_a(t) \\ \dot{q}_b(t) \end{bmatrix} = M \begin{bmatrix} p(t, a) \\ -p(t, b) \end{bmatrix} + \frac{1}{l} M \begin{bmatrix} u(t) \\ -u(t) \end{bmatrix}. \quad (17)$$

Considering equation (5)-(6) together with (13) we deduce

$$p(t, a) = h(t, a^-) - \mu \frac{\partial q}{\partial x}(t, a^-) - h_S(t) + \mu \frac{q_b - q_a}{b - a}, \quad (18)$$

and

$$p(t, b) = h(t, b^+) - \mu \frac{\partial q}{\partial x}(t, b^+) - h_S(t) + \mu \frac{q_b - q_a}{b - a}. \quad (19)$$

Let X be defined by (9), set

$$W := \mathbb{C} \times H^1(\mathcal{E}) \times H^2(\mathcal{E}) \times \mathbb{C} \times \mathbb{C},$$

and denote by $z := [h_S \ h \ q \ q_a \ q_b]^T$ a generic element of X . Consider the operator $A : \mathcal{D}(A) \rightarrow X$ defined by

$$\mathcal{D}(A) := \{z \in W \mid q(a) = q_a, q(b) = q_b\}, \quad (20)$$

$$Az := \begin{bmatrix} -\frac{q(b)-q(a)}{b-a} \\ -\frac{dq}{dx} \\ -\frac{dh}{dx} + \mu \frac{d^2q}{dx^2} \\ R_1z \\ R_2z \end{bmatrix}, \quad (21)$$

where

$$\begin{bmatrix} R_1z \\ R_2z \end{bmatrix} := M \begin{bmatrix} h(a^-) - \mu \frac{dq}{dx}(a^-) - h_S + \mu \frac{q_b - q_a}{b-a} \\ -h(b^+) - \mu \frac{dq}{dx}(b^+) - h_S + \mu \frac{q_b - q_a}{b-a} \end{bmatrix}.$$

We set

$$Bu := \left[0, 0, 0, \frac{lu}{2 \left(1 + \frac{l^3}{12} \right)}, -\frac{lu}{2 \left(1 + \frac{l^3}{12} \right)} \right]^T \text{ and } Cz := h_S, \quad (22)$$

and we observe that $B \in \mathcal{L}(\mathbb{C}, X)$ and $C \in \mathcal{L}(X, \mathbb{C})$.

We have the following result where \mathbb{C}_0 denotes the open right-half plane

$$\mathbb{C}_0 := \{s \in \mathbb{C} : \operatorname{Re} s > 0\}. \quad (23)$$

Proposition

The resolvent set $\rho(A)$ contains \mathbb{C}_0 .

Proposition

The transfer function of the system (1)-(8) is given by  Conflex

$$G(s) := \frac{1}{\left(1 + \frac{l^3}{12}\right) s^2 + \frac{l^2}{2} s \sqrt{1 + \mu s} + \mu l s + l} \quad (s \in \mathbb{C}_0). \quad (24)$$

Lemma

Let F be the function defined by

$$F(s) = \left(1 + \frac{l^3}{12}\right) s^2 + \frac{l^2}{2} s \sqrt{1 + \mu s} + \mu l s + l, \quad (25)$$

and let \mathbb{C}_0 be the open right-half plane, as defined in (23). Then there exists a neighborhood \mathcal{O} of $\overline{\mathbb{C}_0}$ such that F is holomorphic on \mathcal{O} . Moreover, F does not vanish on $\overline{\mathbb{C}_0}$.

Asymptotic behaviour

In this work we study the correct version of this model for vertical displacements of a floating structure, which now reads:



$$\left(1 + \frac{(b-a)^3}{12}\right) \ddot{h}_S(t) = -\frac{(b-a)^2}{2} F * \dot{h}_S(t) - \mu(b-a) \dot{h}_S(t) - (b-a) h_S(t), \quad (26)$$

where μ is the viscosity coefficient of the fluid, $(b-a)$ is the width of the interval $\mathcal{I} = [a, b]$ obtained by projecting the floating object (supposed symmetric around the axis $x = \frac{1}{2}(a+b)$) on the flat horizontal bottom, and $\mathcal{E} = \mathbb{R} \setminus [a, b]$ denotes the viscous fluid domain. Moreover, F is the causal distribution with Laplace transform $\widehat{F}(s) = \sqrt{1 + \mu s}$.

Diffusive representation



Consider the original system, set $\dot{h} := v$ and $z := F * \dot{h}_S$, it can then be viewed as a coupled system

$$\begin{cases} \left(1 + \frac{l^3}{12}\right) \ddot{h}_S + z(t) + \dot{h}_S + l\mu h_S = 0 \\ v(t) = \dot{h}_S(t) \\ \partial_t \varphi(t, \xi) = -\xi \varphi(t, \xi) + v(t); \varphi(0, \xi) = 0 \\ z(t) = \int_{\mu^{-1}}^{\infty} g(\xi) \partial_t \varphi(t, \xi) d\xi + 1v(t). \end{cases} \quad (27)$$

where

$$g(\xi) := \frac{1}{\pi} \frac{\sqrt{\mu\xi - 1}}{\mu\xi}, \quad \text{for } \xi > 1/\mu. \quad (28)$$

Proposition

For all $(h_{S,0}, \omega_0) \in \mathbb{C}^2$, the solution of the coupled system (27), with initial condition $(h_{S,0}, \omega_0, 0)$, satisfies

$$(h_S, \dot{h}_S, \varphi)(t) \rightarrow_{t \rightarrow \infty} 0 \quad \text{in } \mathbb{C}^2 \times \tilde{H},$$

where

$$\tilde{H} = \left\{ \varphi \in L_{\text{loc}}^2(\mathbb{R}^+, d\mathbf{g}), \int_0^\infty \xi |\varphi|^2 d\mathbf{g}(\xi) < \infty \right\}.$$

Proposition

If all the poles s_k of the transfer function lie in the left halfplane $\Re(s) < -\frac{1}{\mu}$, then the asymptotic behaviour of the solution h_S of the system (27) reads

$$h_S(t) \sim K e^{-\frac{t}{\mu}} t^{-3/2}, \quad \text{as } t \rightarrow +\infty.$$

Case $\mu = 0$:

$$h_S(t) = (C_1 \cos(\omega_d t) + C_2 \sin(\omega_d t)) e^{-\delta t}, \quad (29)$$

when $B^2 < 4AI$, where

$$\delta = \frac{B}{2A}, \quad \omega_0 = \sqrt{\frac{I}{A}}, \quad \omega_d = \sqrt{\omega_0^2 - \delta^2} = \frac{\sqrt{4AI - B^2}}{2A}, \quad (30)$$

are the damping coefficient, the undamped natural angular frequency and the damped angular frequency, respectively. The constants C_1 and C_2 , are given by

$$C_1 = h_0, \quad C_2 = \frac{\dot{h}_0 + h_0 \delta}{\omega_d} = \frac{h_0 B + 2\dot{h}_0 A}{\sqrt{4AI - B^2}}. \quad (31)$$

Theorem

The solution of the GFDE is given by

$$h_S(t) = \exp(-\varepsilon t) \left(\sum_{i=1}^4 \Theta_i \mathcal{E}_{\frac{1}{2}}(\lambda_i, t) \right), \quad (32)$$

with constants $\Theta_i := r_i h_0 + \dot{r}_i \dot{h}_0$.

Theorem (Matignon 1996)

We have the following asymptotic equivalents for $\mathcal{E}_\alpha(\lambda, t)$ as t reaches $+\infty$:

- ▶ for $|\arg(\lambda)| \leq \alpha \frac{\pi}{2}$,

$$\mathcal{E}_\alpha(\lambda, t) \sim \frac{1}{\alpha} \lambda^{\frac{1}{\alpha}-1} e^{\lambda^{\frac{1}{\alpha}} t}, \quad (33)$$

- ▶ for $|\arg(\lambda)| > \alpha \frac{\pi}{2}$,

$$\mathcal{E}_\alpha(\lambda, t) \sim \frac{\alpha}{\Gamma(1-\alpha)} \lambda^{-2} t^{-1-\alpha}. \quad (34)$$

- ▶ The case of Proposition 4 is recovered as a special case, which occurs if and only if *all* the roots λ_i fulfill $|\arg(\lambda_i)| > \frac{\pi}{4}$.
- ▶ Otherwise, if *but one* λ_0 lies in the sector $|\arg(\lambda)| < \frac{\pi}{4}$, then a very different asymptotic behaviour is to be found, namely a purely exponentially decaying one, with decay rate $\delta := \varepsilon - \Re(\lambda^2) > 0$ (it must be positive indeed, since asymptotic stability has already been proved in Proposition 3).

Theorem

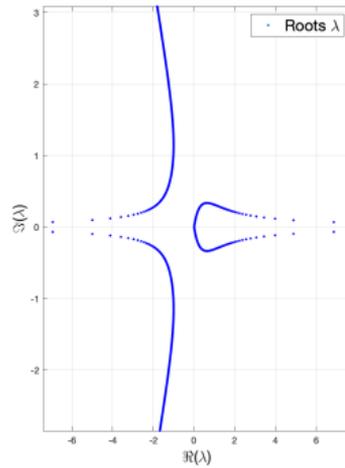


- ▶ *If there is at least one root with $\Re(\lambda_j) > |\Im(\lambda_j)|$ then the asymptotics is of exponential type, with rate $\delta(\mu) := \frac{1}{\mu} - \Re(\lambda^2) > 0$*

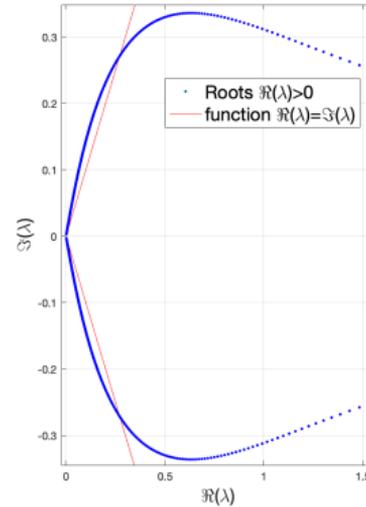
$$h_S(t) \sim \sum_j C_j \exp\left(\left(\lambda_j^2 - \frac{1}{\mu}\right) t\right), \quad (35)$$

- ▶ *or all the four roots lie in $|\arg(\lambda)| > \frac{\pi}{4}$, then the asymptotics is of mixed type,*

$$h_S(t) \sim C t^{-\frac{3}{2}} \exp\left(-\frac{1}{\mu} t\right). \quad (36)$$



(a)



(b)

Figure: Evolution of the four roots λ_i in the σ -plane, as a function of μ . (a): global picture with 4 trajectories. (b): zoom in the right-half plane $\Re(\sigma) > 0$, 2 trajectories crossing the segment $\Re(\lambda) = |\Im(\lambda)|$ for a critical value μ^c of the viscosity.

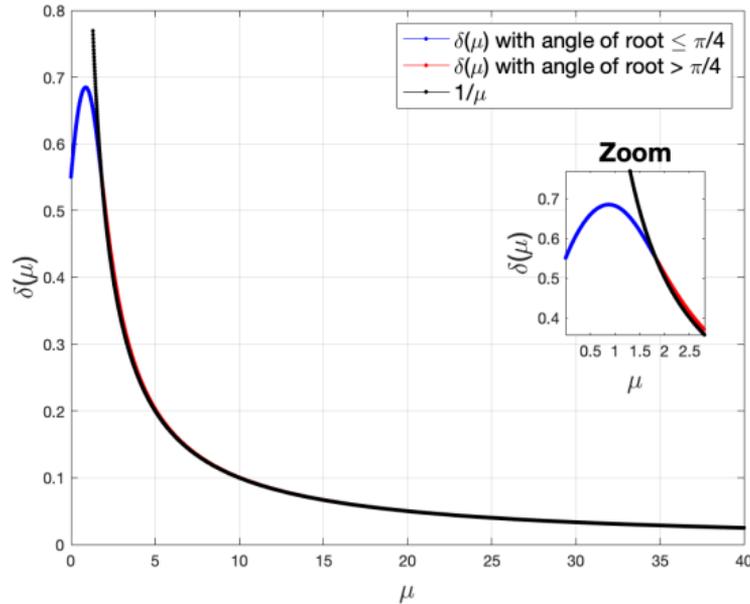


Figure: Damping rate $\delta(\mu) = \Re(\lambda^2) - \frac{1}{\mu}$ as a function of viscosity μ

Thank you !

October 12

Day of indigenous resistance



Figure: Lautaro, leader of the Mapuche resistance.