

# The fractional Brezis–Nirenberg problem in low dimensions

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# The Brezis–Nirenberg problem

For  $N \geq 3$ ,  $\Omega \subset \mathbb{R}^N$  open bounded,  $a \in C(\overline{\Omega})$  [Brezis–Nirenberg 1983] consider

$$\begin{aligned} -\Delta u + au &= N(N-2)u^{\frac{N+2}{N-2}} && \text{in } \Omega, \\ u &> 0 && \text{in } \Omega, \\ u &= 0 && \text{on } \partial\Omega. \end{aligned} \tag{1}$$

[BN] observe: When  $a \equiv 0$  and  $\Omega$  is strictly starshaped, (1) has no solution.

## Proof.

By **Pohozaev's identity** (integrate (1) against  $\nabla u(x) \cdot x$ ), when  $a \equiv 0$

$$-\int_{\partial\Omega} \left| \frac{\partial u}{\partial \nu} \right|^2 \nu(x) \cdot x \, dx = \frac{N-2}{2} \int_{\Omega} |\nabla u|^2 - \frac{N(N-2)^2}{2} \int_{\Omega} u^{\frac{2N}{N-2}} = 0,$$

because  $N(N-2) \int_{\Omega} u^{\frac{2N}{N-2}} = \int_{\Omega} |\nabla u|^2$  (integrate (1) against  $u$ ). If  $\Omega$  is strictly starshaped wrt 0, this implies  $\frac{\partial u}{\partial \nu} \equiv 0$  on  $\partial\Omega$ . Hence  $\int_{\Omega} u^{\frac{N+2}{N-2}} = \int_{\Omega} -\Delta u = -\int_{\partial\Omega} \frac{\partial u}{\partial \nu} = 0$ . But this is a contradiction to  $u > 0$ .  $\square$

So [BN] ask: **Under what conditions on  $a$  does (1) have a solution?**

# Existence of solutions in high and low dimensions

Solutions are given by positive minimizers to the variational problem

$$S(a) := \inf_{u \in H_0^1(\Omega) \setminus \{0\}} \mathcal{S}_a[u], \quad \mathcal{S}_a[u] := \frac{\int_{\Omega} (|\nabla u|^2 + au^2) dx}{\left(\int_{\Omega} |u|^{\frac{2N}{N-2}} dx\right)^{\frac{N-2}{N}}} \quad (2)$$

## Theorem ([Brezis–Nirenberg 1983])

If  $S(a) < S(0)$ , then a minimizer exists (due to E. H. Lieb).

- Let  $N \geq 4$ . Then  $S(a) < S(0) \Leftrightarrow \{a < 0\} \neq \emptyset$ .
- Let  $N = 3$ . Then  $S(a) = S(0)$  whenever  $\|a\|_{\infty}$  is small enough!

**Where does the different behavior in low dimension  $N = 3$  come from ?**

Let  $U_{x,\lambda}(y) = \left(\frac{\lambda}{1+\lambda^2|x-y|^2}\right)^{\frac{N-2}{2}}$  with  $x \in \mathbb{R}^N$  and  $\lambda > 0$ . We notice

$$\|\nabla U_{x,\lambda}\|_{L^2(\mathbb{R}^N)}^2 = S(0) \|U_{x,\lambda}\|_{L^{\frac{2N}{N-2}}(\mathbb{R}^N)}^2 \quad \text{and} \quad -\Delta U_{x,\lambda} = N(N-2)U_{x,\lambda}^{\frac{N+2}{N-2}}.$$

**Idea:** Fix  $x \in \Omega$  and  $\varphi \in C_0^\infty(\Omega)$ . Then test  $\mathcal{S}_a$  with  $\varphi U_{x,\lambda} \in H_0^1(\Omega)$  and let  $\lambda \rightarrow \infty$ !

# The Brezis–Nirenberg dimensional effect

**High dimension:** If  $N \geq 4$ , [BN] compute that, as  $\lambda \rightarrow \infty$ ,

$$\|\nabla(\varphi U_{x,\lambda})\|_2^2 = K_1 + \mathcal{O}(\lambda^{-N+2}),$$

$$\|\varphi U_{x,\lambda}\|_{\frac{2N}{N-2}}^2 = K_2 + \mathcal{O}(\lambda^{-N}),$$

$$\int_{\Omega} a(\varphi U_{x,\lambda})^2 = K_3 a(x) \lambda^{-2} + \mathcal{O}(\lambda^{-N+2}) \quad (K_3 a(x) \lambda^{-2} \log \lambda \text{ if } N = 4).$$

$\implies \mathcal{S}_a[\varphi U_{x,\lambda}] = S(0) + a(x) \lambda^{-2} + \mathcal{O}(\lambda^{-N+2}) < S(0)$  as soon as  $a(x) < 0$  !

**Low dimension:** Let  $N = 3$ ,  $\Omega = B$  and  $a, \varphi$  radial. Then, as  $\lambda \rightarrow \infty$ ,

$$\|\nabla(\varphi U_{x,\lambda})\|_2^2 = K_1 + \lambda^{-1} |\mathbb{S}^2| \int_0^1 |\varphi'(r)|^2 dr + \mathcal{O}(\lambda^{-2}),$$

$$\|\varphi U_{x,\lambda}\|_6^2 = K_2 + \mathcal{O}(\lambda^{-2}),$$

$$\int_{\Omega} a(\varphi U_{x,\lambda})^2 = \lambda^{-1} |\mathbb{S}^2| \int_0^1 a(r) |\varphi(r)|^2 dr + \mathcal{O}(\lambda^{-2}).$$

and hence

$$\mathcal{S}_a[\varphi U_{x,\lambda}] = S(0) + |\mathbb{S}^2| K_2^{-1} \left( \int_0^1 |\varphi'(r)|^2 dr + \int_0^1 a(r) |\varphi(r)|^2 dr \right) \lambda^{-1} + \mathcal{O}(\lambda^{-2}).$$

**Competing lower-order terms!**

# Critical functions and their Green's function

**Question [Brezis 1986]:** Is  $S(a) < S(0)$  also necessary for  $S(a)$  to be achieved?

## Theorem ([Druet 2002])

Let  $N = 3$ . Then  $S(a)$  is achieved  $\Leftrightarrow S(a) < S(0) \Leftrightarrow \{\phi_a < 0\} \neq \emptyset$ .  
As a consequence,  $a$  is critical if and only if  $\min_{\Omega} \phi_a = 0$ .

Here, we define the **Robin function**

$$\phi_a(x) = H_a(x, x), \quad \text{where} \quad G_a(x, y) = \frac{1}{|x - y|^{N-2}} - H_a(x, y)$$

is the Green's function of  $-\Delta + a$ .

Following [Hebey–Vaugon 2001], we call a function  $a \in C(\overline{\Omega})$  **critical** if  $S(a) = S(0)$  and  $S(\tilde{a}) < S(a)$  for all  $\tilde{a} \leq a$  with  $\tilde{a} \neq a$ .

## Corollary

- If  $N \geq 4$ , the only critical function is  $a \equiv 0$ .
- If  $N = 3$ ,  $a$  is critical if and only if  $\min_{\Omega} \phi_a = 0$ . In particular, there are critical functions of all shapes.

# Concentration of solutions

Let  $a$  be critical and let  $(u_\epsilon) \subset H_0^1(\Omega)$  be a sequence of positive solutions to

$$-\Delta u_\epsilon + (a + \epsilon V)u_\epsilon = N(N-2)u_\epsilon^{\frac{N+2}{N-2}} \quad \text{on } \Omega, \quad u_\epsilon = 0 \quad \text{on } \partial\Omega. \quad (3)$$

which blows up at a single point. To leading order,  $u_\epsilon \sim U_{x_\epsilon, \lambda_\epsilon}$  with  $x_\epsilon \rightarrow x_0 \in \bar{\Omega}$  and  $\lambda_\epsilon \rightarrow \infty$ .

In these terms,  $x_0$  is the concentration point and  $\lambda_\epsilon$  is the concentration speed.

Notice that  $\|u_\epsilon\|_\infty \sim \lambda_\epsilon^{\frac{N-2}{2}}$ .

## Theorem

- If  $N \geq 4$  (hence  $a \equiv 0$ ),  $V \equiv -1$ , then  $\nabla \phi_0(x_0) = 0$ . Moreover,  $\lim_{\epsilon \rightarrow 0} \epsilon \|u_\epsilon\|_\infty^{\frac{N-4}{N-2}} = d_N \phi_0(x_0)$ . [Han 1991, Rey 1989]
- If  $N = 3$ , then  $\phi_a(x_0) = 0$ . If  $x_0$  is non-degenerate as a minimum of  $\phi_a$ , with  $a(x_0) < 0$ , then  $\lim_{\epsilon \rightarrow 0} \epsilon \|u_\epsilon\|_\infty^2 = 4\pi^2 \frac{|a(x_0)|}{|\int_\Omega V(y) G_a^2(x_0, y) dy|}$ . [Frank, K., Kovarik CVPDE + arXiv 2021]

# Concentration asymptotics – Proof ideas

**Key idea:** Refine the [BN] test function choice  $\varphi U_{x,\lambda}$  and expand energy quantities to sufficient precision!

We optimize the  $H_0^1$  cutoff procedure: Let  $PU_{x,\lambda}$  be the unique function s.t.

$$-\Delta PU_{x,\lambda} = -\Delta U_{x,\lambda} \text{ in } \Omega, \quad PU_{x,\lambda} = 0 \text{ on } \partial\Omega.$$

**High dimensions:** [Rey 1989]

Let  $N \geq 4$  and  $V \equiv -1$ . As  $\epsilon \rightarrow 0$ , write  $u_\epsilon = PU_{x_\epsilon, \lambda_\epsilon} + w_\epsilon$  with  $\|\nabla w_\epsilon\|_2 = o(1)$  and  $w \perp \{PU_{x,\lambda}, \partial_\lambda PU_{x,\lambda}, \partial_{x_i} PU_{x,\lambda}\}$ . Expand Pohozaev

$$\epsilon \int_{\Omega} u_\epsilon^2 = \int_{\partial\Omega} (x - x_0) \cdot \nu(x) \left| \frac{\partial u_\epsilon}{\partial \nu} \right|^2$$

using the quantitative bound  $\|w_\epsilon\|_2 = \mathcal{O}(\lambda_\epsilon^{-\frac{N+2}{2}})$ , we get

$$a_N \epsilon \lambda_\epsilon^2 + o(\epsilon \lambda_\epsilon^2) = b_N \phi_0(x_0) \lambda_\epsilon^{-N+2} + o(\lambda_\epsilon^{-N+2}).$$

**Low dimension:** [Frank, K., Kovarik 2021]

When  $N = 3$ , terms in energy expansion are again arranged differently. To expand to the desired precision, we need the refined development

$$u_\epsilon = PU_{x_\epsilon, \lambda_\epsilon} + \lambda_\epsilon^{-1/2} (H_0(x_\epsilon, \cdot) - H_a(x_\epsilon, \cdot)) + q_\epsilon, \quad \|\nabla q_\epsilon\|_2^2 = \lambda_\epsilon^{-1} + \epsilon \lambda_\epsilon^{-1/2}.$$

## Going fractional

The appropriate generalization of the BN problem to fractional orders of derivatives  $2s < N$  with  $s \in (0, 1)$  is

$$\begin{aligned} (-\Delta)^s u + (a + \epsilon V)u &= c_{N,s} u^{\frac{N+2s}{N-2s}} && \text{in } \Omega, \\ u &> 0 && \text{in } \Omega, \\ u &= 0 && \text{on } \mathbb{R}^N \setminus \Omega. \end{aligned} \tag{4}$$

with the **fractional Laplacian**  $(-\Delta)^s$  given as

$$(-\Delta)^s u = \mathcal{F}^{-1}(|\xi|^{2s} \mathcal{F}u) = C_{N,s} P.V. \int_{\mathbb{R}^N} \frac{u(x) - u(y)}{|x - y|^{N+2s}} dy.$$

The variational problem associated to (4) is

$$S_{N,s}(a) := \inf \frac{\int_{\mathbb{R}^N} |(-\Delta)^{s/2} u|^2 + \int_{\Omega} a u^2}{\left( \int_{\Omega} u^{\frac{2N}{N-2s}} \right)^{\frac{N-2s}{N}}}$$

where the inf is taken over  $\tilde{H}^s(\Omega) := \{u \in H^s(\mathbb{R}^N) : u \equiv 0 \text{ on } \mathbb{R}^N \setminus \Omega\}$ .

The notions of Green's function  $G_a(x, y) = \frac{1}{|x - y|^{N-2s}} - H_a(x, y)$ , and critical function  $a$  carry over without problem.



## Some known results for the fractional BN problem

The analysis of the fractional BN problem with  $s \in (0, 1)$  presents some additional difficulties which mostly stem from the fact that the operator  $(-\Delta)^s$  is **non-local**.

**Still:**  $U_{x,\lambda}(y) = \left( \frac{\lambda}{1+\lambda^2|x-y|^2} \right)^{\frac{N-2s}{2}}$  satisfies  $(-\Delta)^s U_{x,\lambda} = c_{N,s} U_{x,\lambda}^{\frac{N+2s}{N-2s}}$ .

### Theorem (Servadei–Valdinoci 2013, 2015)

- If  $N \geq 4s$ , then  $S_{N,s}(a) < S_{N,s}$  whenever  $a(x) < 0$  for some  $x \in \Omega$ .
- If  $2s < N < 4s$ , then  $S_{N,s}(a) < S_{N,s}$  if  $a(x) < -\mu_s < 0$ .

(Compare [Brezis–Nirenberg 1983].)

### Theorem (Choi–Kim–Lee 2014)

If  $N > 4s$ , and solutions  $u_\epsilon$  to (4) with  $a \equiv 0$ ,  $V = -1$  blow up at exactly one point  $x_0 \in \Omega$ , then  $\lim_{\epsilon \rightarrow 0} \epsilon \|u_\epsilon\|_\infty^{2\frac{n-4s}{n-2s}} = \delta_{N,s} \phi_0(x_0)$ . (Compare [Han 1991, Rey 1989].)

$\Rightarrow$  **What about low dimensions**  $2s < N < 4s$  ??

# Main results

## Theorem 1 (De Nitti, K., 2021)

Let  $s \in (0, 1)$  and  $2s < N < 4s$  and let  $a \in C(\overline{\Omega})$ . Then

$$(i) S_{N,s}(a) \text{ is achieved} \Leftrightarrow (ii) S_{N,s}(a) < S_{N,s} \Leftrightarrow (iii) \{\phi_a < 0\} \neq \emptyset.$$

As a consequence,  $a$  is critical if and only if  $\min_{\Omega} \phi_a = 0$ .

- This is the fractional version of [Druet 2002].
- $(i) \Leftrightarrow (ii)$  follows similarly to [BN].  $(ii) \Leftrightarrow (iii)$  follows from

$$S_a[\psi_{x,\lambda}] = S(0) + c\phi_a(x)\lambda^{-N+2s} + o(\lambda^{-N+2s}) \quad (5)$$

as  $\lambda \rightarrow \infty$ , where  $\psi_{x,\lambda} = PU_{x,\lambda} + \lambda^{-\frac{N-2s}{2}}(H_0(x, \cdot) - H_a(x, \cdot))$ .

- The forward implications are the hard part. E.g.  $(iii) \Rightarrow (ii)$  requires to show: For  $a$  critical and  $u_\epsilon$  minimizers of  $S(a - \epsilon)$ , one has

$$S(0) > S(a - \epsilon) = S(0) + c(\phi_a(x_\epsilon) + o(1))\lambda^{-N+2s}.$$

Thus  $\phi_{a-\epsilon}(x_0) < \phi_a(x_0) \leq 0$  and  $(iii)$  follows.

# Main results

Write  $Q_V(x) := \int_{\Omega} V(y) G_a(x, y)^2 dy$ .

## Theorem 2 (De Nitti, K., 2021)

Let  $s \in (0, 1)$  and  $\frac{8}{3}s < N < 4s$  and let  $a \in C(\overline{\Omega})$  be critical. Suppose that  $(u_{\epsilon}) \subset \tilde{H}^s(\Omega)$  is a sequence of energy-minimizing solutions to (4). Then the  $u_{\epsilon}$  blow up in exactly one point  $x_0 \in \Omega$  satisfying  $\phi_a(x_0) = 0$ . Moreover,

$$u_{\epsilon} = PU_{x_{\epsilon}, \lambda_{\epsilon}} + \lambda_{\epsilon}^{-\frac{N-2s}{2}} (H_0(x_{\epsilon}, \cdot) - H_a(x_{\epsilon}, \cdot)) + q_{\epsilon}$$

with  $\|((-\Delta)^{s/2})q_{\epsilon}\|_2^2 = o(\lambda_{\epsilon}^{-2s} + \lambda_{\epsilon}^{-2N+4s})$  and  $x_{\epsilon} \rightarrow x_0$  and

$$\lim_{\epsilon \rightarrow 0} \epsilon \lambda_{\epsilon}^{4s-N} = d_{N,s} \frac{|a(x_0)|}{|Q_V(x_0)|}. \quad (6)$$

The point  $x_0$  maximizes  $\frac{|Q_V(x_0)|^{\frac{2s}{4s-N}}}{|a(x_0)|^{\frac{N-2s}{4s-N}}}$  among all  $x$  with  $\phi_a(x) = 0$  and  $V(x) < 0$ .

This is the fractional version of [\[Frank, K., Kovarik CVPDE 2021\]](#)

## Some comments

- Our proof strategy (for both Theorems 1 and 2) is **variational**. In particular Theorem 2 also holds for almost-minimizers  $u_\epsilon$  of  $S_{N,s}(a + \epsilon V)$ , which need not satisfy any PDE.
- If we knew that  $\|q_\epsilon\|_\infty = o(\lambda_\epsilon^{\frac{N-2s}{2}})$ , Theorem 2 would yield the value of  $\lim_{\epsilon \rightarrow 0} \epsilon \|u_\epsilon\|_\infty^{2\frac{4s-N}{N-2s}}$ .
- For  $s = 1$ , the recent preprint [Frank, K., Kovarik arXiv 2021] removes the energy-minimizing assumption, thus proving a conjecture from [Brezis–Peletier 1989]. The analogous question for  $s \in (0, 1)$  is open.
- For  $N > 4s$  we prove (work in preparation)

$$\lim_{\epsilon \rightarrow 0} \epsilon \lambda_\epsilon^{N-4s} = d_{N,s} \frac{\phi_0(x_0)}{|V(x_0)|}.$$

and  $x_0$  achieves  $\max_{\{x: V(x) < 0\}} \phi_0(x)^{-\frac{2s}{N-4s}} |V(x)|^{\frac{N-2s}{N-4s}}$

- If  $N < 4s$ , the 'renormalized energy'  $|Q_V(x)|^{\frac{2s}{4s-N}} |a(x)|^{-\frac{N-2s}{4s-N}}$  is non-local in  $V$ !

# The proof for fractional $s \in (0, 1)$

## Basic results:

- Concentration compactness [Palatucci–Pisante 2015] (Compare [Struwe 1984].)
- $PU_{x,\lambda}$  and orthogonality conditions [Abdelhedi–Chtioui–Hajajiej 2017] (Compare [Bahri–Coron 1988].)

Thus we can write  $u_\epsilon = \alpha_\epsilon(PU_{x_\epsilon, \lambda_\epsilon} + w_\epsilon)$  with  $\alpha_\epsilon \rightarrow 1$ ,  $\|(-\Delta)^s w_\epsilon\|_{L^2(\mathbb{R}^N)} = o(1)$  and  $w_\epsilon \perp \{PU_{x,\lambda}, \partial_\lambda PU_{x,\lambda}, \partial_{x_i} PU_{x,\lambda}\}$ .

## New ingredients:

- Precise analysis of the functions  $PU_{x,\lambda}$  and  $H_a(x, y)$
- Spectral **coercivity inequality**: For  $p = \frac{2N}{N-2s}$ ,

$$\|(-\Delta)^{s/2} w_\epsilon\|_2^2 - c_{N,s}(p-1) \int_{\Omega} U_{x_\epsilon, \lambda_\epsilon}^{p-2} w_\epsilon^2 \geq \frac{4s}{N+2s+2} \|(-\Delta)^{s/2} w_\epsilon\|_2^2.$$

by stereographic projection to  $\mathbb{S}^N$ . (Compare [Rey 1990, Bianchi–Egnell 1991].  
For  $s \in (0, 1)$ , see also [Chen–Frank–Weth 2012]. )

- Non-existence of a minimizer for  $S(a)$  when  $a$  critical. (Compare [Druet 2002].)

## The proof for fractional $s \in (0, 1)$ - cont'd

- The assumption  $8s/3 < N$  should be technical. The phenomenon that makes it necessary to be imposed has no analogue when  $s = 1$  and  $N = 3$ : To absorb error terms in

$$S(a + \epsilon V) = S_{N,s} + (\phi_a(x) + \epsilon \int_{\Omega} VG_a(x_{\epsilon}, \cdot))\lambda^{-N+2s} + a(x)\lambda^{-2s} \\ + o(\lambda^{-2s}) + (\epsilon + \phi_a(x))\lambda^{-N+2s}.$$

we need  $\lambda^{-k(N-2s)} \stackrel{!}{=} o(\lambda^{-2s})$  for all  $k \geq 3$ .

It would be very interesting to understand more precisely the impact of such 'lower-order BN dimensional effects'!

- Simplifications / differences with respect to previous works
  - ▶ Avoid the intermediate spectral cutoff argument from [\[Frank–K.–Kovarik CVPDE 2021\]](#).
  - ▶ Avoid the formulation of  $(-\Delta)^s$  as a Dirichlet-to-Neumann problem for a degenerate-elliptic local PDE in  $N + 1$  dimension as in [\[Choi–Kim–Lee 2014\]](#) and others

# Summary

- **Brezis–Nirenberg problem:** On  $\Omega \subset \mathbb{R}^N$ , positive solutions to  $(-\Delta)^s u + au = u^{\frac{N+2s}{N-2s}}$  exist sometimes, but not always.
  - ▶ The picture is more intricate in low dimensions  $2s < N < 4s$  !
- **Main Result:** In  $2s < N < 4s$ :
  - ▶  $a$  is critical  $\Leftrightarrow \min_{\Omega} \phi_a = 0$ .
  - ▶ Let  $(u_{\epsilon})$  almost minimizers of the perturbed critical Brezis–Nirenberg energy  $S(a + \epsilon V)$  at  $a$  critical. We precisely characterize
    - ★ the concentration point  $x_0$  and the concentration speed  $\lambda_{\epsilon}$  in terms of the data  $a, V, \Omega$
    - ★ the blow-up profile:  $u_{\epsilon} \sim PU_{x_{\epsilon}, \lambda_{\epsilon}} + \lambda_{\epsilon}^{-\frac{N-2s}{2}} (H_0(x_{\epsilon}, \cdot) - H_a(x_{\epsilon}, \cdot))$ .
- **Method of proof:** Purely variational  $\rightarrow$  Refine iteratively the energy expansion and improve error bounds by exploiting coercivity + energy minimality.
- **Perspectives:**
  - ▶ Treat very low dimensions  $2s < N \leq \frac{8s}{3}$  by understanding better the related cancellation phenomenon.
  - ▶ Remove energy bounds: simple blow-up for non-minimizers, multibubble blow-up.

Vielen Dank für Eure Aufmerksamkeit! :-)