

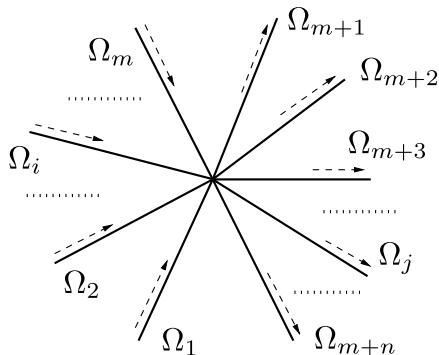
Conservation laws on a star-shaped network

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FAU DCN-AvH Seminar, 2 March 2023

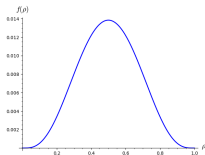
We consider a junction consisting of m incoming and n outgoing edges.



- Incoming edges: $x \in \Omega_i = \mathbb{R}_-, i = 1, \dots, m$;
- Outgoing edges: $x \in \Omega_j = \mathbb{R}_+, j = m + 1, \dots, m + n$;
- The junction is located at $x = 0$.

On each edge we consider the evolution problem

$$\partial_t \rho_h + \partial_x f_h(\rho_h) = 0, \quad h = 1, \dots, m+n,$$



- ρ_h conserved quantity,
- f_h flux : possibly different, non degenerate nonlinear and bell-shaped
 - $f_h : [0, R] \rightarrow \mathbb{R}_+$, Lipschitz continuous,
 - $f_h(0) = 0 = f_h(R)$,
 - $\exists \bar{\rho} \in [0, R]$, such that $f'_h(\rho)(\bar{\rho} - \rho) > 0$, for a.e. $\rho \in [0, R]$.

We postulate **conservation** at the junction

$$\frac{d}{dt} \sum_{h=1}^{m+n} \int_{\Omega_h} \rho_h(t, x) dx = 0,$$

which we rewrite as

$$\sum_{i=1}^m f_i(\rho_i(t, 0^-)) = \sum_{j=m+1}^{m+n} f_j(\rho_j(t, 0^+)).$$

Weak solutions, edge-wise entropy admissible

We call **weak solution** on the star-shaped network $\vec{\rho} = (\rho_1, \dots, \rho_{m+n})$

- $\rho_h \in L^\infty(\mathbb{R}_+ \times \Omega_h; [0, R])$;
- ρ_h is a Kruzhkov entropy solution in $\mathbb{R}_+ \times \{\Omega_h \setminus \partial\Omega_h\}$.
Namely $\forall k \in [0, R]$ and $\forall \varphi \in \mathcal{C}_c^1(\mathbb{R}_+ \times \Omega_h)$, $\varphi \geq 0$

$$\int_{\mathbb{R}_+} \int_{\Omega_h} |\rho_h - k| \varphi_t + \text{sign}(\rho_h - k) (f_h(\rho_h) - f_h(k)) \varphi_x \, dx \, dt \\ + \int_{\Omega_h} |u_0^h(x) - k| \varphi(0, x) \, dx \geq 0 ;$$

- **conservation** at the junction holds.

➔ **Weak solutions** are not unique in general.

Different approaches to solve the Riemann problem at the junction

(edge-wise constant initial conditions)

- Holden-Risebro 1995
maximize a concave "entropy" function at the junction ;
- Coclite-Piccoli 2002, Coclite-Garavello-Piccoli 2005
traffic distribution matrix + optimization ;
- Lebacque 1996, Lebacque-Khoshyaran 2002
Supply-Demand model.
Link between junctions and boundary value problems ;
- Coclite-Garavello 2010, Andreianov-Coclite-D. 2017, Coclite-D. 2020
Vanishing Viscosity approximation, Germ of admissible solutions
- Fjordholm-Musch-Risebro 2022 (2 papers)
Finite volumes approximation, Germ of admissible solutions ;
- Holle 2021, (Holle-Herty-Westdickenberg 2020, p-systems)
Kinetic BGK approximation with maximal entropy dissipation .

The junction as a family of IBVPs

Fix $\vec{u}_0 = (u_0^1, \dots, u_0^{m+n})$

We look for $\vec{\rho} = (\rho_1, \dots, \rho_{m+n})$ s.t. $\forall h, \rho_h \in L^\infty(\mathbb{R}_+ \times \Omega_h, [0, R])$ solves

$$\begin{cases} \partial_t \rho_h + \partial_x f_h(\rho_h) = 0, & \text{on }]0, T[\times \Omega_h, \\ \rho_h(t, 0) = v_h(t), & \text{on }]0, T[, \\ \rho_h(0, x) = u_0^h(x), & \text{on } \Omega_h, \end{cases}$$

where $\vec{v} : \mathbb{R}_+ \rightarrow [0, R]^{m+n}$ is to be fixed at each $t > 0$

- ▶▶ to ensure conservation,
- ▶▶ depending on the state of the system,
- ▶▶ depending on the model.

... Solves? Weak entropy solution for the IBVP

u is a weak entropy solution for the IBVP

$$\begin{cases} \partial_t u + \partial_x f(u) = 0, & \text{for } (t, x) \text{ in } \mathbb{R}_+ \times \mathbb{R}_- \\ u(t, 0^-) = u_b(t), \\ u(0, x) = u_0(x), \end{cases}$$

if

- u is a Kruzhkov entropy solution in the interior of $\mathbb{R}_+ \times \mathbb{R}_-$,
- u satisfies the boundary condition in the sense of Bardos-LeRoux-Nédélec

$$\begin{aligned} \text{sign}(u(t, 0^-) - u_b(t))(f(u(t, 0^-)) - f(k)) &\geq 0, \\ &\forall k \in \mathcal{I}(u(t, 0^-), u_b(t)), \end{aligned}$$

which also write as

$$f(u(t, 0^-)) = \text{God}(u(t, 0^-), u_b(t)).$$

Geometrical meaning of the (BLN) condition

The solution of

$$\begin{cases} \partial_t u + \partial_x f(u) = 0, & \text{for } (t, x) \text{ in } \mathbb{R}_+ \times \mathbb{R}_- \\ u(t, 0^-) = u_b, \\ u(0, x) = u_0, \end{cases}$$

is the restriction to $\mathbb{R}_+ \times \mathbb{R}_-$ of the Riemann problem

$$\begin{cases} \partial_t u + \partial_x f(u) = 0, & \text{for } (t, x) \text{ in } \mathbb{R}_+ \times \mathbb{R} \\ u(0, x) = \begin{cases} u_0, & x \leq 0, \\ u_b, & x > 0. \end{cases} \end{cases}$$

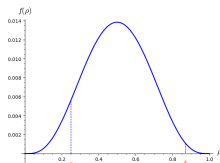
Then we say that the boundary condition is satisfied as soon as

$$\begin{cases} \partial_t u + \partial_x f(u) = 0, & \text{for } (t, x) \text{ in } \mathbb{R}_+ \times \mathbb{R} \\ u(0, x) = \begin{cases} u(t, 0^-), & x \leq 0, \\ u_b, & x > 0. \end{cases} \end{cases}$$

only contains waves of positive speed.

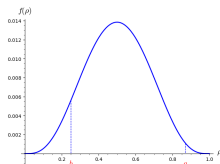
Godunov's flux

$$\begin{cases} \partial_t u + \partial_x f(u) = 0, & (t, x) \in \mathbb{R}_+ \times \mathbb{R}, \\ u_0(x) = \begin{cases} a, & \text{if } x < 0, \\ b, & \text{if } x > 0. \end{cases} \end{cases}$$



The Godunov flux is the function $(a, b) \mapsto f(u(t, 0^-)) = f(u(t, 0^+))$. Analytically

$$\text{God}(a, b) = \begin{cases} \min_{s \in [a, b]} f(s) & \text{if } a \leq b, \\ \max_{s \in [b, a]} f(s) & \text{if } a \geq b. \end{cases}$$



Basic properties:

- Consistency: for all $a \in [0, R]$, $\text{God}(a, a) = f(a)$;
- Monotonicity and Lipschitz continuity: $\exists L > 0$ s.t. $\forall (a, b) \in [0, R]^2$

$$0 \leq \partial_a \text{God}(a, b) \leq L, \quad -L \leq \partial_b \text{God}(a, b) \leq 0.$$

The Coclite-Garavello-Piccoli Riemann solver

To describe the Riemann Solver at the junction we define

▶▶ for $i = 1, \dots, m$

Demand function : $\Delta_i(\rho_i(t, 0^-)) = \max_s \text{God}_{f_i}(\rho_i(t, 0^-), s)$;

▶▶ for $j = m + 1, \dots, m + n$

Supply function : $\Sigma_j(\rho_j(t, 0^+)) = \max_s \text{God}_{f_j}(s, \rho_j(t, 0^+))$;

and we use them to determine the **passing flow** at the junction from each of the incoming roads

$$\Gamma_i : [0, R]^3 \rightarrow [0, f_i^{\max}], \quad i = 1, \dots, m.$$

At a 1-2 divide

Fix a **distribution factor** $\beta \in (0, 1)$.

The **passing flow** at the junction is $\Gamma_1 : [0, R]^3 \rightarrow [0, f_1^{\max}]$ such that :

- ▶ If $\beta \Delta_1(\rho_{1,0}) \leq \Sigma_2(\rho_{2,0})$, and $(1 - \beta) \Delta_1(\rho_{1,0}) \leq \Sigma_3(\rho_{3,0})$ then $\Gamma_1(\vec{\rho}_0) = \Delta_1(\rho_{1,0})$,
- ▶ otherwise, $\Gamma_1(\vec{\rho}_0) = \min\{\beta^{-1} \Sigma_2(\rho_{2,0}), (1 - \beta)^{-1} \Sigma_3(\rho_{3,0})\}$.

In both cases

$$\left\{ \begin{array}{l} v_1 = \left(f_1|_{[\bar{\rho}_1, R]} \right)^{-1} (\Gamma_1), \\ v_2 = \left(f_2|_{[0, \bar{\rho}_2]} \right)^{-1} (\beta \Gamma_1), \\ v_3 = \left(f_3|_{[0, \bar{\rho}_3]} \right)^{-1} ((1 - \beta) \Gamma_1). \end{array} \right.$$

Remark

The application $\vec{\rho}_0 \mapsto \vec{v}$ is not monotone :

$$\rho_{h,0} \geq \tilde{\rho}_{h,0} \not\Rightarrow v_h \geq \tilde{v}_h.$$

At a 2 – 2 junction

We introduce a distribution matrix of the form

$$A = \begin{pmatrix} \beta & \gamma \\ 1 - \beta & 1 - \gamma \end{pmatrix}$$

with β and γ in $]0, 1[\setminus \{1/2\}$.

Then

- $(\Gamma_1, \Gamma_2) \in [0, \Delta_1] \times [0, \Delta_2]$;
- $A \cdot (\Gamma_1, \Gamma_2)^T$ must be in $[0, \Sigma_3] \times [0, \Sigma_4]$;
- $\Gamma_1 + \Gamma_2$ should be as large as possible, under the constraints above.

Remark

Counterexamples show that this solver lacks L^1 -Lipschitz continuity with respect to the initial conditions.

See [Coclite-Garavello-Piccoli, 2005] and the book by Garavello and Piccoli *Traffic Flow on Networks*.

Vanishing viscosity approximations

[Coclite-Garavello, 2010]

Fix $\varepsilon > 0$ and consider

$$\begin{cases} \partial_t \rho_h^\varepsilon + \partial_x f_h(\rho_h^\varepsilon) = \varepsilon \partial_{xx}^2 \rho_h^\varepsilon, \\ \sum_{i=1}^m (f_i(\rho_i^\varepsilon(t, 0)) - \varepsilon \partial_x \rho_i^\varepsilon(t, 0)) = \sum_{j=m+1}^{m+n} (f_j(\rho_j^\varepsilon(t, 0)) - \varepsilon \partial_x \rho_j^\varepsilon(t, 0)), \\ \rho_h^\varepsilon(t, 0) = \rho_{h'}^\varepsilon(t, 0), \\ \rho_h^\varepsilon(0, x) = u_{h,\varepsilon}^0(x), \end{cases}$$

where the approximated initial conditions $\vec{u}_{0,\varepsilon}$ satisfy

$$u_{h,\varepsilon}^0 \in W^{2,1} \cap C^\infty(\Omega_h; [0, R]),$$

$$u_{h,\varepsilon}^0 \rightarrow \rho_{0,h}, \text{ a.e. and in } L^p(\Omega_h), \quad 1 \leq p < \infty, \text{ as } \varepsilon \rightarrow 0,$$

$$\|u_{h,\varepsilon}^0\|_{L^1(\Omega_h)} \leq \|\rho_{0,h}\|_{L^1(\Omega_h)}, \quad \|\partial_x u_{h,\varepsilon}^0\|_{L^1(\Omega_h)} \leq TV(\rho_{0,h}), \quad \varepsilon \|\partial_{xx}^2 u_{h,\varepsilon}^0\|_{L^1(\Omega_h)} \leq C_0,$$

with $C_0 > 0$ independent from ε , h .

Theory of semigroups \Rightarrow for any fixed $\varepsilon > 0$ there exists a unique ρ^ε s.t.

$$\rho_h^\varepsilon \in C([0, \infty); L^2(\Omega_h)) \cap L^1_{loc}((0, \infty); W^{2,1}(\Omega_h)), \quad \forall h,$$

$$0 \leq \rho_h^\varepsilon \leq R, \quad \sum_{h=1}^{m+n} \|\rho_h^\varepsilon(t, \cdot)\|_{L^1(\Omega_h)} \leq \sum_{h=1}^{m+n} \|\rho_{0,h}\|_{L^1(\Omega_h)}, \quad \forall t \geq 0,$$

+ additional a priori estimates.

Compensated compactness \Rightarrow existence of a sequence $\{\varepsilon_\ell\}_{\ell \in \mathbb{N}}$, $\varepsilon_\ell \rightarrow 0$ and a **weak solution** $\bar{\rho}$ of the inviscid Cauchy problem at the junction s.t.

$$\rho_h^{\varepsilon_\ell} \longrightarrow \rho_h, \text{ a.e. and in } L^p_{loc}(\mathbb{R}_+ \times \Omega_h), \quad 1 \leq p < \infty,$$

for every $h \in \{1, \dots, m+n\}$.

In [Andreianov-D.-Coclite, 2017] we further characterize the limit solution and prove its uniqueness. More details in the following. . .

Analogously to [Diehl, 2009], [Andreianov-Mitrović, 2015] for $m = n = 1$

The condition $\rho_h^\varepsilon(t, 0) = \rho_{h'}^\varepsilon(t, 0)$, $\forall h, h' \in \{1, \dots, m+n\}$, translates into

$$v_h(t) = v_{h'}(t),$$

for the family of hyperbolic IBVPs at the junction.

$\vec{\rho} = (\rho_1, \dots, \rho_{m+n})$ is an **admissible solution** if there exists v in $L^\infty(\mathbb{R}_+, [0, R])$ s.t.

- $\vec{\rho}$ is a **weak solution**,
- each component ρ_h is weak entropy solution for the IBVP

$$\begin{cases} \rho_{h,t} + f_h(\rho_h)_x = 0, & \text{on }]0, T[\times \Omega_h, \\ \rho_h(t, 0) = v(t), & \text{on }]0, T[, \\ \rho_h(0, x) = u_0^h(x), & \text{on } \Omega_h. \end{cases}$$

We call **germ of vanishing viscosity** the set

$$\mathcal{G}_{VV} = \left\{ \vec{k} \in [0, R]^{m+n}, \text{ stationary edge-wise constant } \text{admissible solution} \right\}$$

Lemma

If ρ_h is a Kruzhkov entropy solution in the interior of $\mathbb{R}_+ \times \Omega_h$, $\forall h \in \{1, \dots, m+n\}$,
TFAE

- $\vec{\rho}$ is an **admissible solution**;
- for a.e. $t \in \mathbb{R}_+$, the vector of traces $\gamma \vec{\rho}(t) = (\rho_1(t, 0^-), \dots, \rho_{m+n}(t, 0^+))$ is in \mathcal{G}_{VV} ;
- $\forall \vec{k} \in \mathcal{G}_{VV}$, $\vec{\rho}$ satisfies **adapted entropy inequality** on the network:
 $\forall \xi \in \mathcal{D}(\mathbb{R}_+ \times \mathbb{R}), \xi \geq 0$,

$$\sum_{h=1}^{m+n} \left(\int_{\mathbb{R}_+} \int_{\Omega_h} \{ |\rho_h - k_h| \xi_t + \text{sign}(\rho_h - k_h) (f_h(\rho_h) - f_h(k_h)) \xi_x \} dx dt \right) \geq 0.$$

Well-posedness for admissible solutions

Theorem

- For any \vec{u}_0 there exists at most one **admissible solution** $\vec{\rho}$.
- If $\vec{\rho}$ and $\vec{\rho}^*$ are **admissible solutions** corresponding to \vec{u}_0 and \vec{v}_0 , then

$$\sum_{h=1}^{m+n} \|\rho_h(t) - \rho_h^*(t)\|_{L^1(\Omega_h; \mathbb{R})} \leq \sum_{h=1}^{m+n} \left\| u_h^0 - v_h^0 \right\|_{L^1(\Omega_h; \mathbb{R})}$$

Fundamental properties of \mathcal{G}_{VV}

- **completeness** : we can associate an **admissible solution** to any Riemann datum.
- **dissipativity** : for any \vec{k}_1, \vec{k}_2 in \mathcal{G}_{VV} with $\vec{k}_\ell = (k_1^\ell, \dots, k_{m+n}^\ell)$, $\ell = 1, 2$,

$$\sum_{i=1}^m \text{sign}(k_i^1 - k_i^2) \left(f_i(k_i^1) - f_i(k_i^2) \right) - \sum_{j=m+1}^{m+n} \text{sign}(k_j^1 - k_j^2) \left(f_j(k_j^1) - f_j(k_j^2) \right) \geq 0.$$

- **maximality** : if \vec{k}_1 satisfies \uparrow for all \vec{k}_2 in \mathcal{G}_{VV} , then $\vec{k}_1 \in \mathcal{G}_{VV}$.

Vanishing viscosity with different coupling conditions?

See [Guaraguaglini–Natalini, 2015 & 2021] for the linear case

We consider coupling conditions inspired by the Kedem-Katchalsky conditions for membrane permeability

$$\begin{cases} \partial_t \rho_h^\varepsilon + \partial_x f_h(\rho_h^\varepsilon) = \varepsilon \partial_{xx}^2 \rho_h^\varepsilon, & t > 0, x \in \Omega_h, \\ \rho_h^\varepsilon(0, x) = \rho_{h,0}^\varepsilon(x), & h = 1, \dots, m+n, \\ f_i(\rho_i^\varepsilon(t, 0)) - \varepsilon \partial_x \rho_i^\varepsilon(t, 0) = \sum_j c_{ij}(\rho_i^\varepsilon(t, 0) - \rho_j^\varepsilon(t, 0)), & i = 1, \dots, m, \\ f_j(\rho_j^\varepsilon(t, 0)) - \varepsilon \partial_x \rho_j^\varepsilon(t, 0) = \sum_i c_{ij}(\rho_i^\varepsilon(t, 0) - \rho_j^\varepsilon(t, 0)), & j = m+1, \dots, m+n, \end{cases}$$

where $c_{ij} > 0$. We do not impose continuity at $x = 0$.

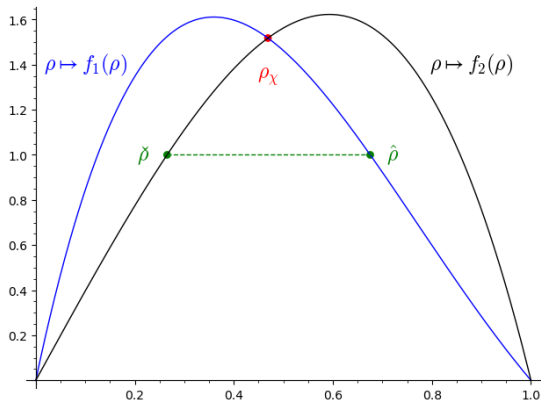
We can prove [Coclite–D. 2020]:

- Existence of parabolic approximations for any ε ;
- Convergence (up to a subsequence) to a **weak solution**.

Characterization of the hyperbolic limit?

The 1-1 case

Assume $f_1(\rho_x) = f_2(\rho_x)$ and $f_1(\hat{\rho}) = f_2(\check{\rho}) = c(\hat{\rho} - \check{\rho})$



$$\begin{cases}
 \partial_t \rho_1^\varepsilon + \partial_x f_1(\rho_1^\varepsilon) = \varepsilon \partial_{xx}^2 \rho_1^\varepsilon, & t > 0, x < 0, \\
 \partial_t \rho_2^\varepsilon + \partial_x f_2(\rho_2^\varepsilon) = \varepsilon \partial_{xx}^2 \rho_2^\varepsilon, & t > 0, x > 0, \\
 f_1(\rho_1^\varepsilon(t, 0)) - \varepsilon \partial_x \rho_1^\varepsilon(t, 0) = c(\rho_1^\varepsilon(t, 0) - \rho_2^\varepsilon(t, 0)), & t > 0, \\
 f_2(\rho_2^\varepsilon(t, 0)) - \varepsilon \partial_x \rho_2^\varepsilon(t, 0) = c(\rho_1^\varepsilon(t, 0) - \rho_2^\varepsilon(t, 0)), & t > 0, \\
 \rho_1^\varepsilon(0, x) = \hat{\rho}, & x < 0, \\
 \rho_2^\varepsilon(0, x) = \check{\rho}, & x > 0.
 \end{cases}$$

As $\varepsilon \rightarrow 0$ the limit is $\rho_1(t, x) \equiv \hat{\rho}$, $\rho_2(t, x) \equiv \check{\rho}$.

- The couple $(\hat{\rho}, \check{\rho})$ is a **connection** as introduced by [Adimurthi-Mishra-Gowda, 2005].
- Already obtained by adapted vanishing viscosity regularization

$$\begin{cases}
 \partial_t \rho_1^\varepsilon + \partial_x f_1(\rho_1^\varepsilon) = \varepsilon \partial_{xx}^2 a_1(\rho_1^\varepsilon), & t > 0, x < 0, \\
 \partial_t \rho_2^\varepsilon + \partial_x f_2(\rho_2^\varepsilon) = \varepsilon \partial_{xx}^2 a_2(\rho_2^\varepsilon), & t > 0, x > 0,
 \end{cases}$$

where a_1 and $a_2 : [0, 1] \rightarrow [0, 1]$ are strictly monotone increasing bijections and $a_1(\hat{\rho}) = a_2(\check{\rho})$.

Can we find a complete, maximal, L^1 dissipative germ for the limit problem?

$\mathcal{G}_{KK} = \{(u_L, u_R) \text{ s.t. } u(t, x) = u_L \mathbb{1}_{\mathbb{R}_-} + u_R \mathbb{1}_{\mathbb{R}_+} \text{ is a stationary admissible solution}\}$

$A = (\hat{\rho}, \check{\rho})$ defined by $f_1(\hat{\rho}) = f_2(\check{\rho}) = c(\hat{\rho} - \check{\rho})$ must be in, together with $u(t, x) \equiv 0$ and $u(t, x) \equiv 1$.

Of course, the germ contains all the couples (a, b) which

- are traces of IBVPs with boundary conditions $\vec{v} = (\hat{\rho}, \check{\rho})$, $\vec{v} = (0, 0)$ or $\vec{v} = (1, 1)$;
- satisfy the Rankine-Hugoniot condition.

$$\mathcal{H} = \left\{ \begin{array}{l} (a, b): a \in [0, \check{u}_L] \cup \{\hat{\rho}\}, \\ b \in [\hat{u}_R, 1] \cup \{\check{\rho}\}, \\ f_1(a) = f_2(b). \end{array} \right\} \cup \{(0, 0)\} \cup \{(1, 1)\},$$

where \check{u}_L and \hat{u}_R satisfy $f_1(\check{u}_L) = f_1(\hat{\rho})$ and $f_2(\hat{u}_R) = f_2(\check{\rho})$

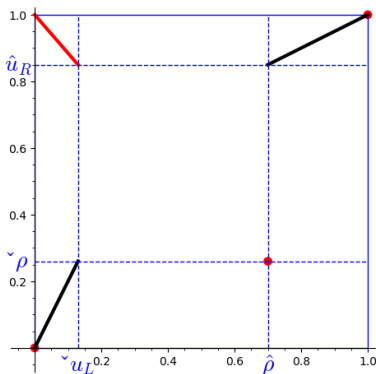
Some Riemann problems do not have solutions if we impose that the traces at $x = 0$ of the solution are in \mathcal{H} .

We call \mathcal{G}_A^* the set of all couples (a, b) such that $f_1(a) = f_2(b)$ and

$$\text{sign}(a - \hat{\rho}) (f_1(a) - f_1(\hat{\rho})) - \text{sign}(b - \check{\rho}) (f_2(b) - f_2(\check{\rho})) \geq 0,$$

This set need to be in the germ because we need maximality and L^1 -dissipativity.

A case by case study show that \mathcal{G}_A^* is complete and maximal. Therefore, we can use it to prove well-posedness of limit solutions to the hyperbolic problem.



Thank you for your attention!