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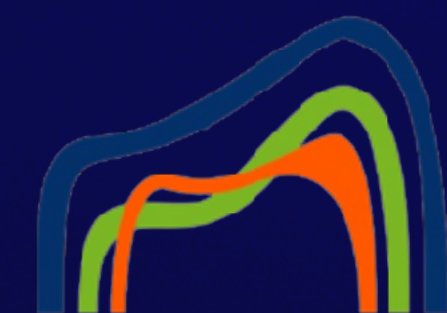
Trends in the Mathematical Sciences

Networked Hyperbolic Systems

Modeling, Control and Efficient Simulation

Yue Wang
June 11, 2024

DFG Deutsche
Forschungsgemeinschaft

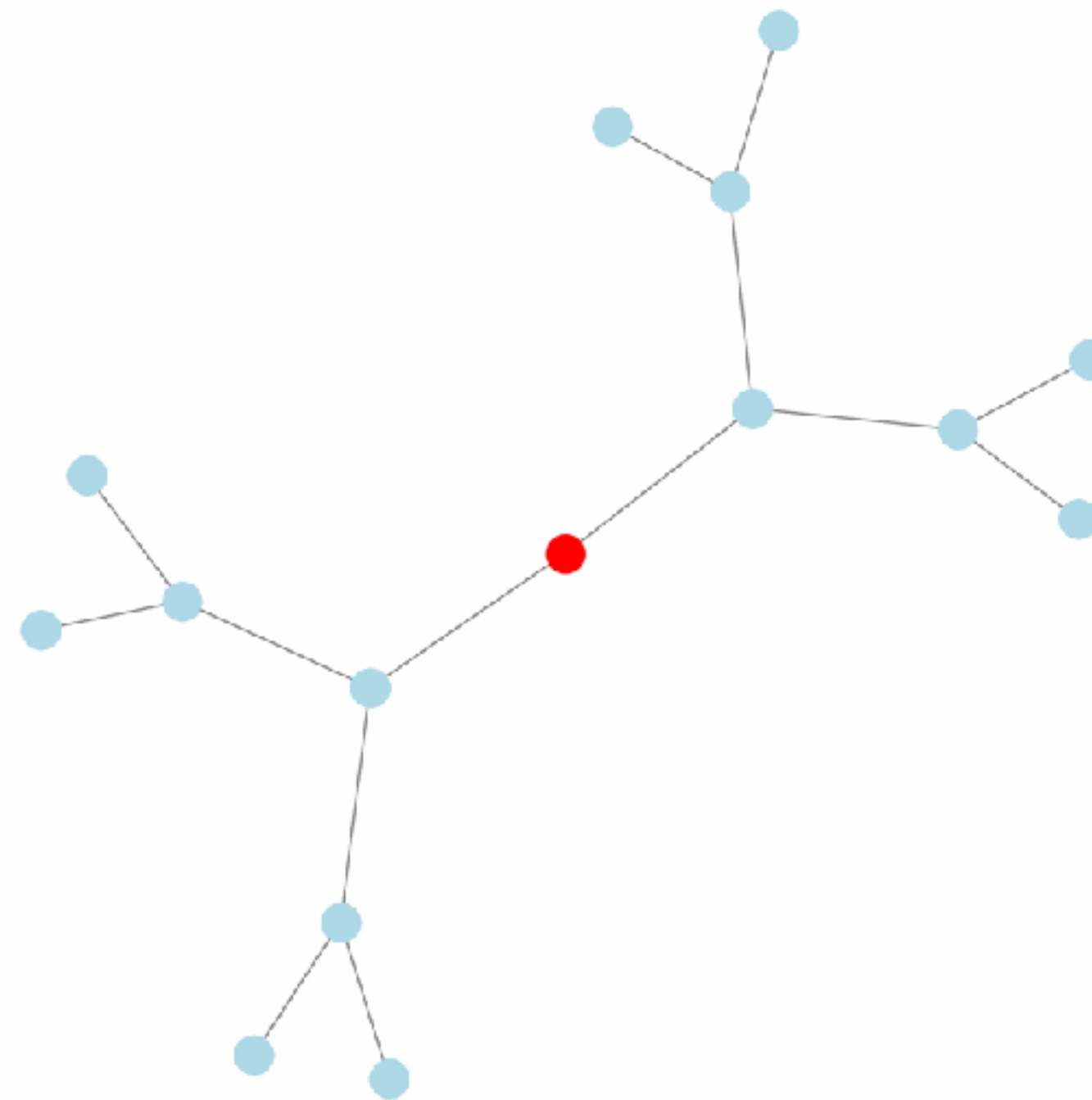


Friedrich-Alexander-Universität
DYNAMICS, CONTROL,
MACHINE LEARNING
AND NUMERICS

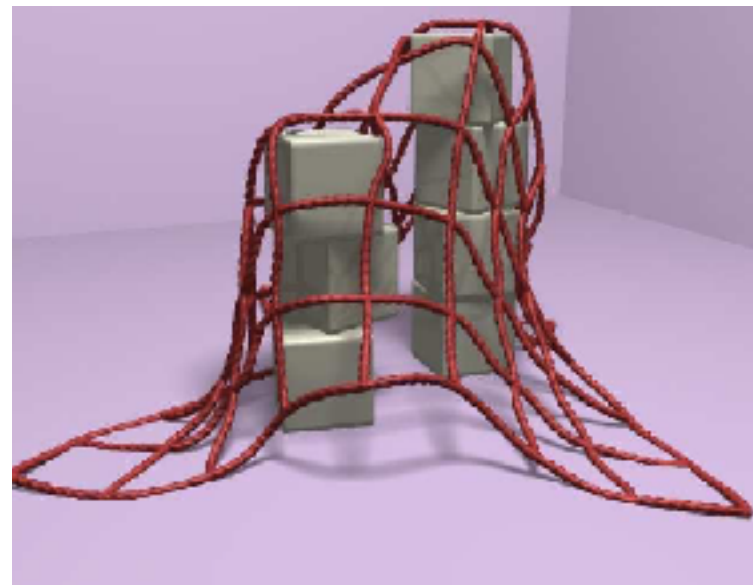
Networked Hyperbolic Systems

Hyperbolic systems are a class of partial differential equations (PDEs) that describe **wave propagation** and other phenomena where **information travels with finite speed**.

When these systems are interconnected in a network, they form networked hyperbolic systems.



Real-world Applications



Network of Large Deflection Strings
(Nonlinear coupled wave equations)



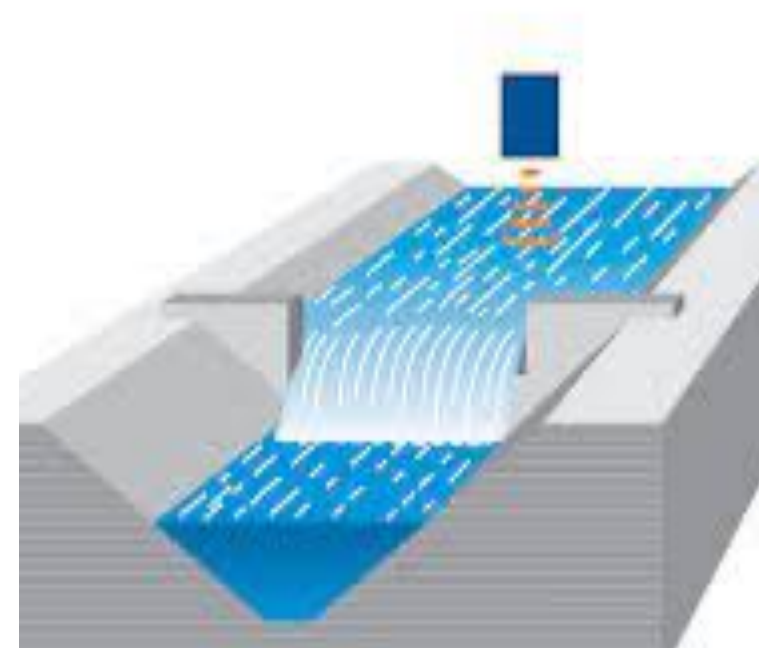
NASA Flexible Flight Device
(Geometrically Exact Beams)



Wind Turbine



Gas transport networks
(Isothermal Euler Equations)



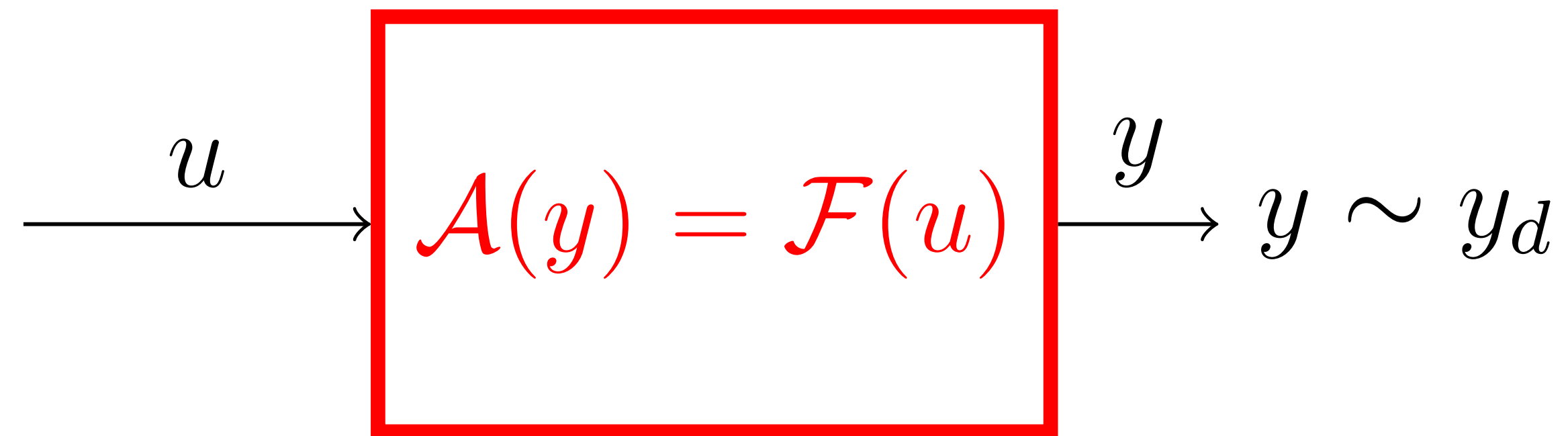
Open Canal
(Saint-Venant Equation)



Flexible Robotic Arm

- Project **ConFlex (2017–2022)** and **ModConFlex (2023–2027)**: Modeling and Control for Flexible Structures Interacting with Fluids
- Project **SFB TRR154 (2018–2026)**: Mathematical modelling, simulation and optimization using the example of gas networks
- Project **DFG WA5144/1-1 (2022–2024)**: Analysis and Control of Nonlinear Hyperbolic Systems with Degeneration on Networks
- Sino-German Mobility Project **CIN-PDE (2022-2025)**: Control, Inversion and Numerics for Partial Differential Equations (Coming workshop in October 8th. -10th. at Shanghai)

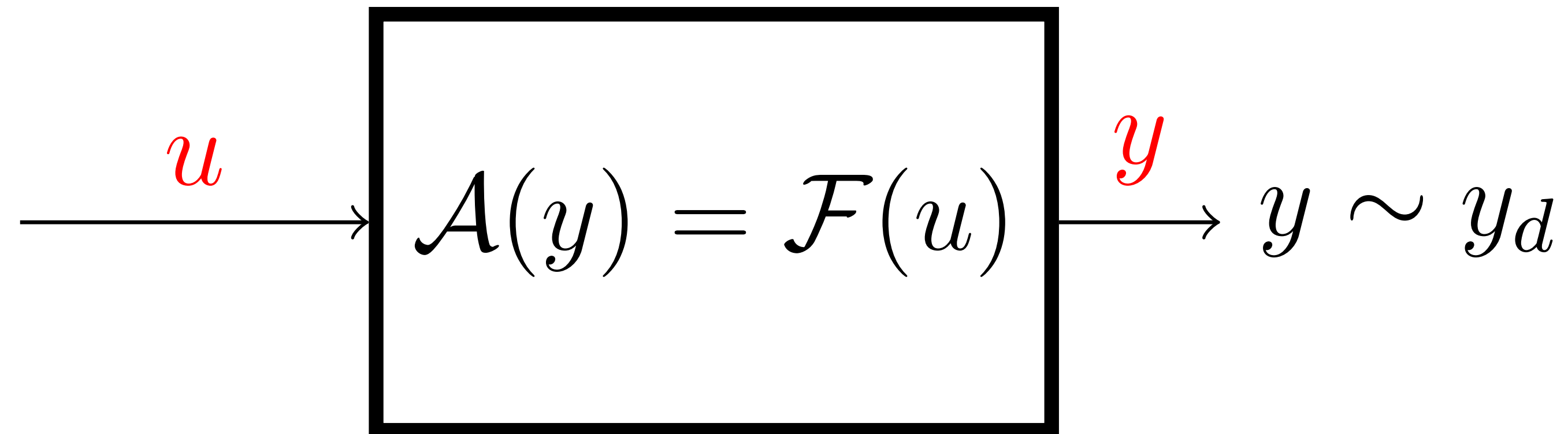
Significant Interests



► **Modeling and Analysis:**

- 1 Physics-driven & Data-driven (e.g. Physic Laws & Machine Learning).
- 2 For analysis, difficulties may arise on networks with
 - **nonlinear** elements,
 - non-trivial **boundary conditions** and **coupling**,
 - complex **topological structure**.

Significant Interests

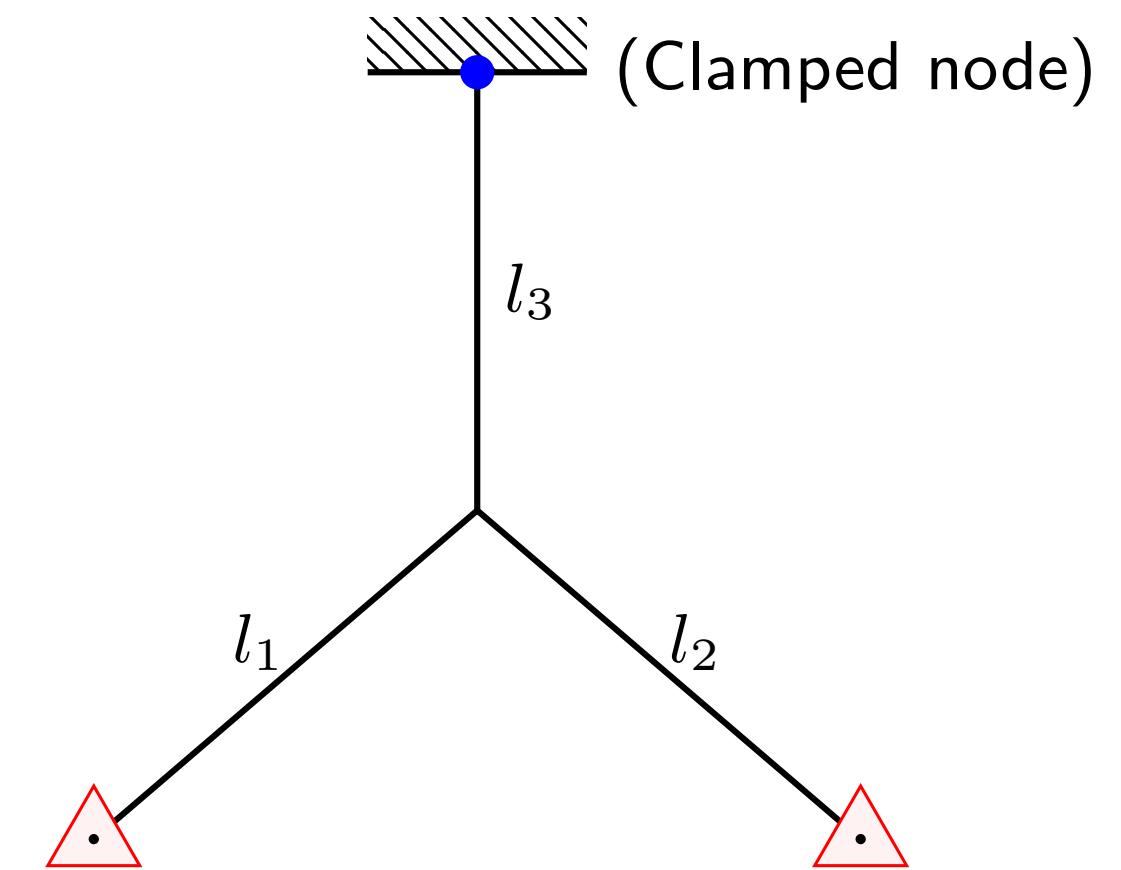


- ▶ Modeling and Analysis
- ▶ Control Theory and Optimal Design
 - 1 Feasibility \longrightarrow Controllability (To find at least one way to reach the target, e.g. $y(x, T) = y_d(x)$ (Exact Controllability), or $y(N_i, t) = y_d(t)$ (Nodal Profile Control));
 - 2 Optimality \longrightarrow Optimal control (To find the best way, in some sense, to reach the target. e.g. controllability time, minimum number of controls, control design on networks, constrained optimization)

Example: Networks of vibrating strings

Two boundary control problems (linear case)

	Null Controllability	Controllability of Nodal Profile
Control Target	$y^i(T, x) = y_t^i(T, x) = 0$	$y_x^3(t, l_3) = y_d, \quad t > T$
(sharp) controllability time $T \geq T^*$	$T^* = 2(l_3 + \max\{l_1, l_2\})$	$T^* = l_3 + \min\{l_1, l_2\}$
minimum number of required controls	2	1



$$\left\{ \begin{array}{l} y_{tt}^i - y_{xx}^i = 0, \quad (t, x) \in (0, T) \times (0, l_i), \quad i = 1, 2, 3, \\ x = 0 : y^1(t, 0) = y^2(t, 0) = y^3(t, 0), \quad t \in (0, T), \\ y_x^0(t, 0) + y_x^1(t, 0) + y_x^2(t, 0) = 0, \quad t \in (0, T), \\ x = l_3 : y^3(t, l_3) = 0, \quad t \in (0, T), \\ x = l_i : y_x^i(t, l_i) = u^i(t), \quad t \in (0, T), \quad i = 1, 2, \end{array} \right.$$

- René Dáger, Enrique Zuazua. *Wave Propagation, Observation and Control in 1-d Flexible Multi-Structures* (2006)
- Yue Wang, Günter Leugering. *Boundary controllability and observability of nodal profile for wave equation* (2022)

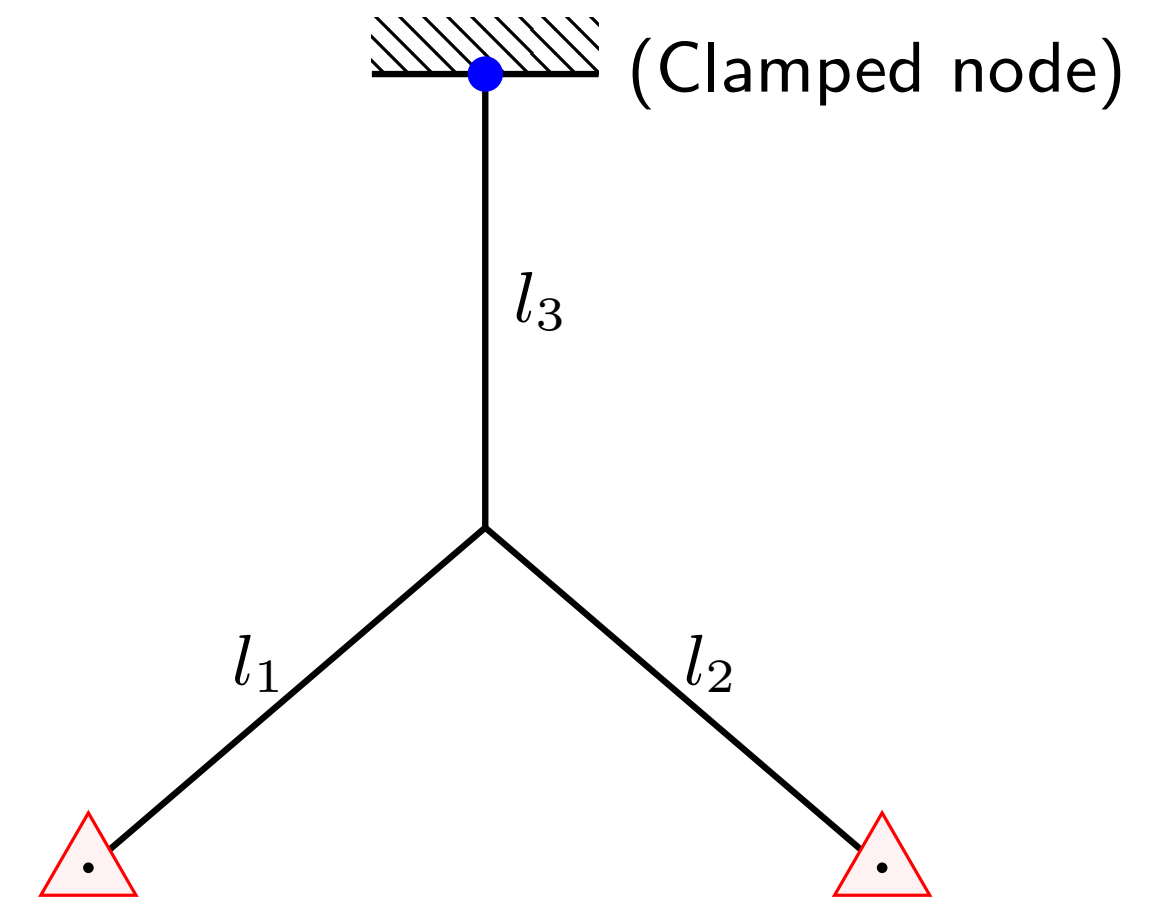
Key Techniques: Set suitable Hilbert space, duality between controllability and observability, observability Inequality.

Example: Networks of vibrating strings

Nonlinear Case

Consider the following coupled system of 1-D quasilinear wave equations ($i = 1, \dots, n$):

$$(\mathbf{E}) \begin{cases} y_{tt}^i - (K^i(y^i, y_x^i))_x = F(\mathbf{y}, \mathbf{y}_x, \mathbf{y}_t), & x \in [0, L_i], t \in [0, T] \\ \sum K^i(y^i(t, 0), y_x^i(t, 0)) = 0, & t \in [0, T] \\ y^j(t, 0) = y^i(t, 0), & i \neq j, \\ y^i(t, L_i) = u^i(t), & t \in [0, T] \\ (y^i, y_t^i)(0, x) = (\phi^i(x), \psi^i(x)), & x \in [0, L_i]. \end{cases}$$



where

- ▶ $\mathbf{y} = (y^1, \dots, y^n)^T$ is an unknown vector function of (t, x) ,
- ▶ $K^i = K^i(y^i, y_x^i)$ are given C^2 functions of y^i and y_x^i ,
- ▶ $\frac{\partial}{\partial y_x^i} K^i(y^i, y_x^i) > 0$,
- ▶ u^i can be considered as 0 (no control) or control function.

!! HUM method (J.Lions, 1980s) and duality method (E.Zuazua, 1990s) can not be applied on this case.

Difficulties may arise in...

- ▶ **Nonlinearity.**

- > **Weak solutions.** [of quasilinear hyperbolic systems \rightarrow shock waves \rightarrow an irreversible process \rightarrow Impossible to get exact boundary controllability for any arbitrarily given initial and final states [A. Bressan, G. M. Coclite, '02]

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- > **Classical solution** exists only locally in time (P. D. Lax, '64; F. John, '90; T. Li, '94) → **semi-global classical solution ($T > 0$ might be suitably large)** [M. Cirinà, '70, T. Li, Y. Jin, B. Rao, '00, '01] → Local exact controllability in the quasilinear case.

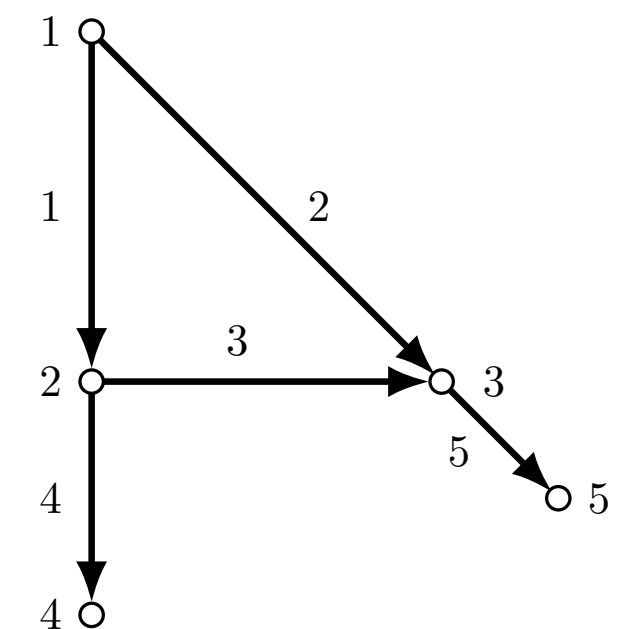
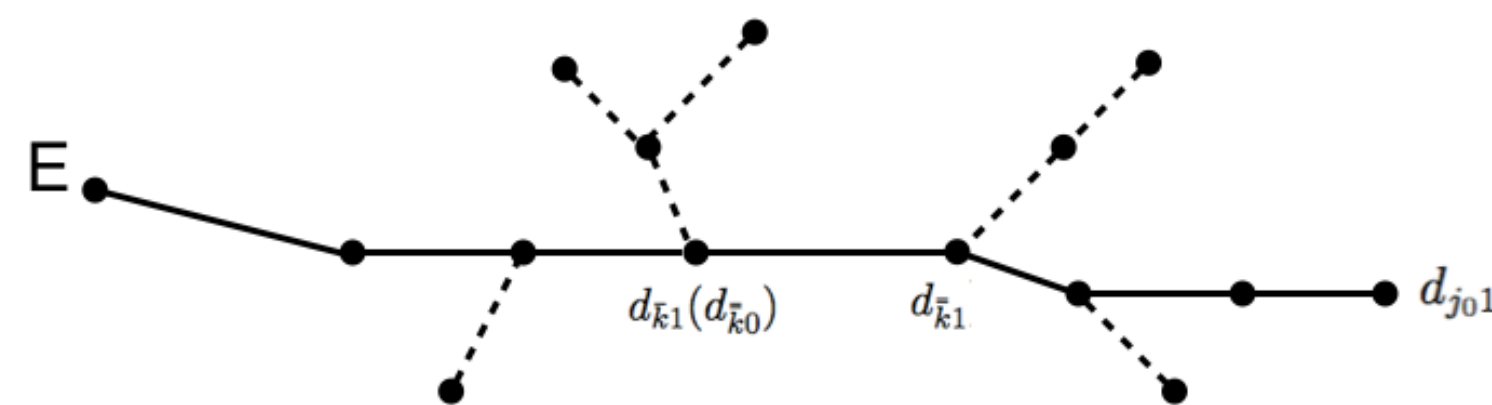
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► Networked Structure.

- > **Coupling at the junction.** Complexity and Nonlinearity in **interface conditions.**
- > Complex topological structure of networks $\mathcal{G} = (\mathcal{V}, \mathcal{E})$ may change the controllability results [Lagnese-Leugeing-Schmidt, '94]



Example: Networks of vibrating strings

New boundary conditions + coupling

Consider the following **coupled** system of 1-D **quasilinear** wave equations ($i = 1, \dots, n$):

$$(\mathbf{E}) \left\{ \begin{array}{l} y_{tt}^i - (K^i(y^i, y_x^i))_x = F(\mathbf{y}, \mathbf{y}_x, \mathbf{y}_t), \quad x \in [0, L_i], t \in [0, T] \\ y_{tt}^i(t, 0) = G^i(t, \mathbf{y}(t, 0), \mathbf{y}_x(t, 0), \mathbf{y}_t(t, 0)) \\ \quad + \int_0^t H^i(t, s, \mathbf{y}(s, 0)) ds, \quad t \in [0, T] \\ y^i(t, L_i) = u^i(t), \quad t \in [0, T] \\ (y^i, y_t^i)(0, x) = (\phi^i(x), \psi^i(x)), \quad x \in [0, L_i]. \end{array} \right.$$

Second-order differential operators

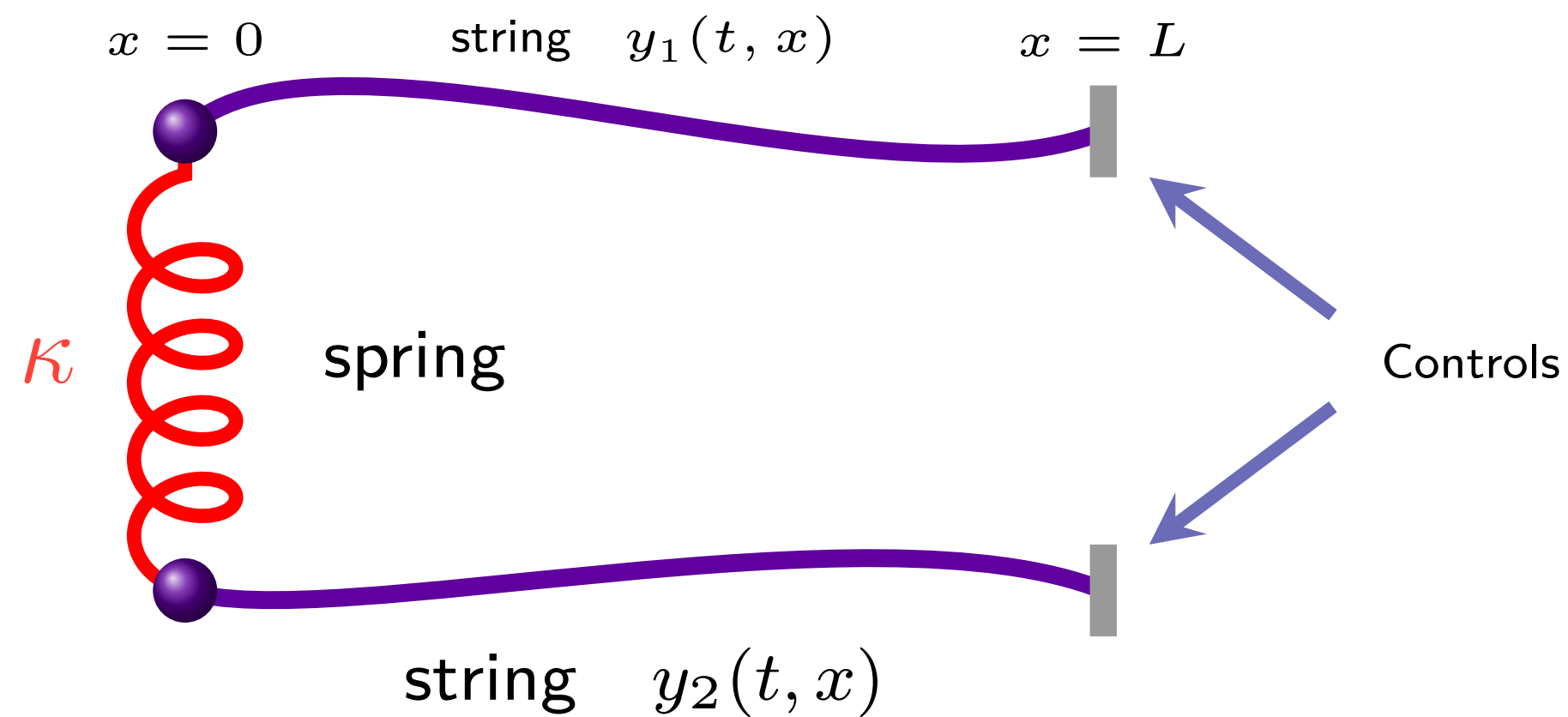
(temporal) non-locality

where

- ▶ $\mathbf{y} = (y^1, \dots, y^n)^T$ is an unknown vector function of (t, x) ,
- ▶ $K^i = K^i(y^i, y_x^i)$ are given C^2 functions of y^i and y_x^i ,
- ▶ $\frac{\partial}{\partial y_x^i} K^i(y^i, y_x^i) > 0$,
- ▶ F^i, G^i, H^i are given C^1 functions of their arguments and 0 value at null state (i.e. 0 is an equilibrium).

Example: String-mass-spring system

$$\left\{ \begin{array}{l} y_{tt}^i - K_i (y_x^i)_x = 0, \quad 0 \leq x \leq L, t > 0, \quad i = 1, 2, \\ x = 0 : y_{tt}^1(0, t) = K_1 (y_x^1(0, t)) - \kappa (y^1(0, t) - y^2(0, t)), \\ \quad y_{tt}^2(0, t) = K_2 (y_x^2(0, t)) + \kappa (y^1(0, t) - y^2(0, t)), \\ x = L : y^i = u^i(t), \quad i = 1, 2. \end{array} \right. \Rightarrow \text{Dynamical transmission conditions}$$



Example: String-mass-spring system

$$\left\{ \begin{array}{l} y_{tt}^i - K_i (y_x^i)_x = 0, \quad 0 \leq x \leq L, t > 0, \quad i = 1, 2, \\ x = 0 : y_{tt}^1(0, t) = K_1 (y_x^1(0, t)) - \kappa (y^1(0, t) - y^2(0, t)), \\ \quad y_{tt}^2(0, t) = K_2 (y_x^2(0, t)) + \kappa (y^1(0, t) - y^2(0, t)), \\ x = L : y^i = u^i(t), \quad i = 1, 2. \end{array} \right.$$

- ▶ If the spring stiffness tends to infinity, formally the system tends to the classical string-mass problem.¹
- ▶ For string-mass system it is known that the mass smoothens the waves while crossing the mass-point.²
- ▶ If the spring stiffness tends to zero, the strings become uncoupled.
- ▶ This system is controllable by only 1 control, and in this case, we discovered **asymmetric solution space** (smoothing effect of mass on waves) and **controllable space**.³

¹G. Leugering, 1998; F. Almusallams, 2015; Y.Wang, T.Li, 2018

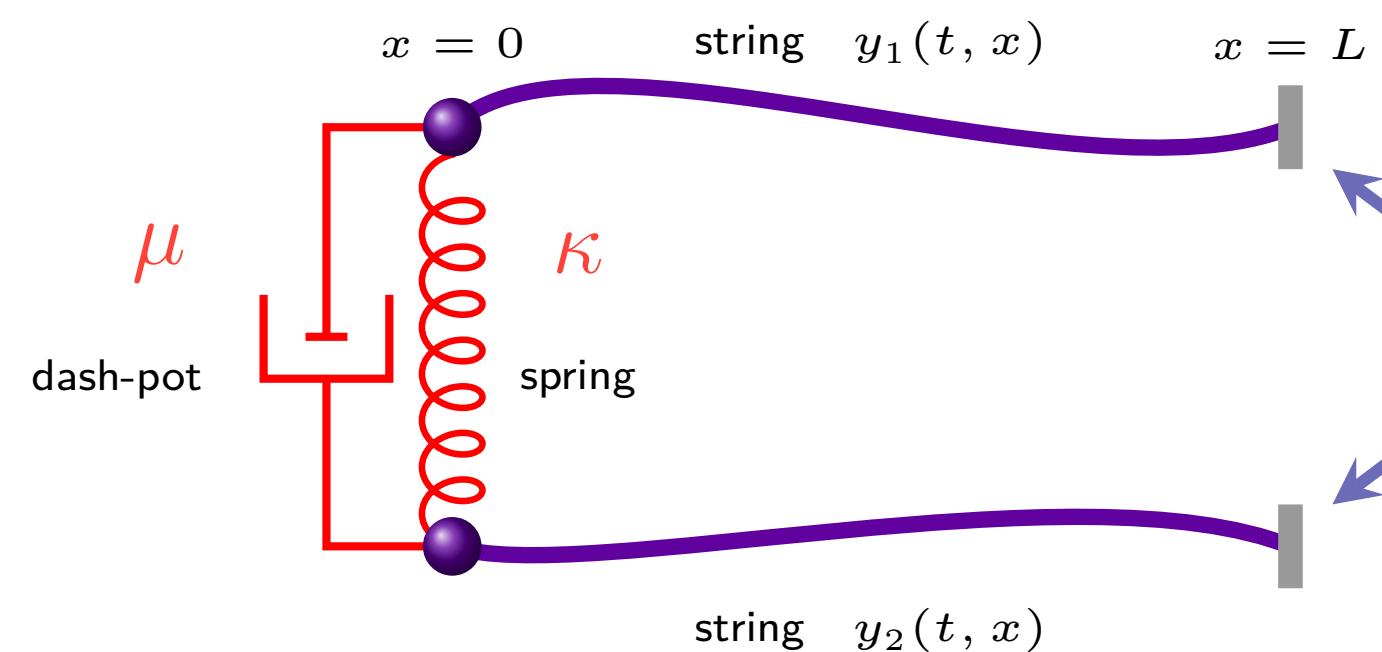
²S. Hansen, E.Zuazua 1995

³G.Leugering,S.Micu, I.Roventa, Y.Wang, 2022

Other examples for dynamical boundary conditions

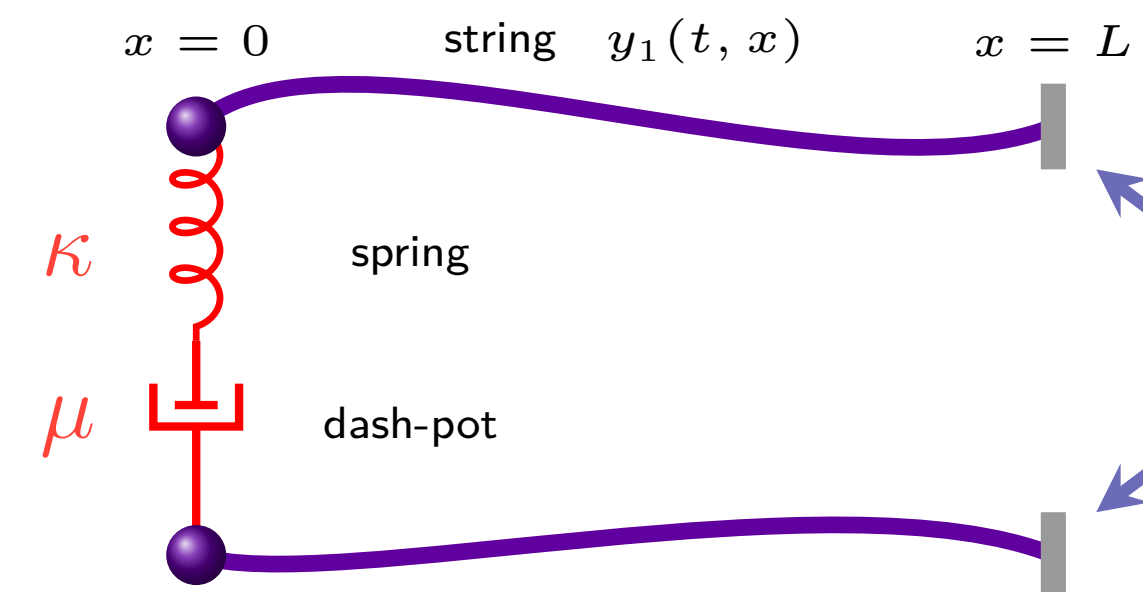
- **Kelvin Model:** a classical class of viscoelastic **solid** models.

$$\begin{cases} y_{tt}^i - K_i(y_x^i)_x = 0, & 0 \leq x \leq L, t > 0, & i = 1, 2, \\ x = 0 : y_{tt}^1(t, 0) = K_1(y_x^1(t, 0)) - \kappa(y^1(t, 0) - y^2(t, 0)) + \mu(y_t^1(t, 0) - y_t^2(t, 0)), \\ & y_{tt}^2(t, 0) = K_2(y_x^2(t, 0)) + \kappa(y^1(t, 0) - y^2(t, 0)) + \mu(y_t^1(t, 0) - y_t^2(t, 0)), \\ x = L : y^i = u^i(t), & i = 1, 2. \end{cases}$$



- **Maxwell Model:** a classical class of viscoelastic **fluid** models.

$$\begin{cases} y_{tt}^i - K_i(y_x^i)_x = 0, & 0 \leq x \leq L, & i = 1, 2, \\ x = 0 : y_{tt}^1(0, t) = K_1(y_x^1(t, 0)) - \kappa(y^1(t, 0) - y^2(t, 0)) \\ & + \frac{\kappa^2}{\mu} \int_0^t e^{-\frac{\kappa}{\mu}(t-\tau)} (y^1(\tau, 0) - y^2(\tau, 0)) d\tau \\ & y_{tt}^2(0, t) = \dots \\ x = L : y^i = u^i(t), & i = 1, 2. \end{cases}$$



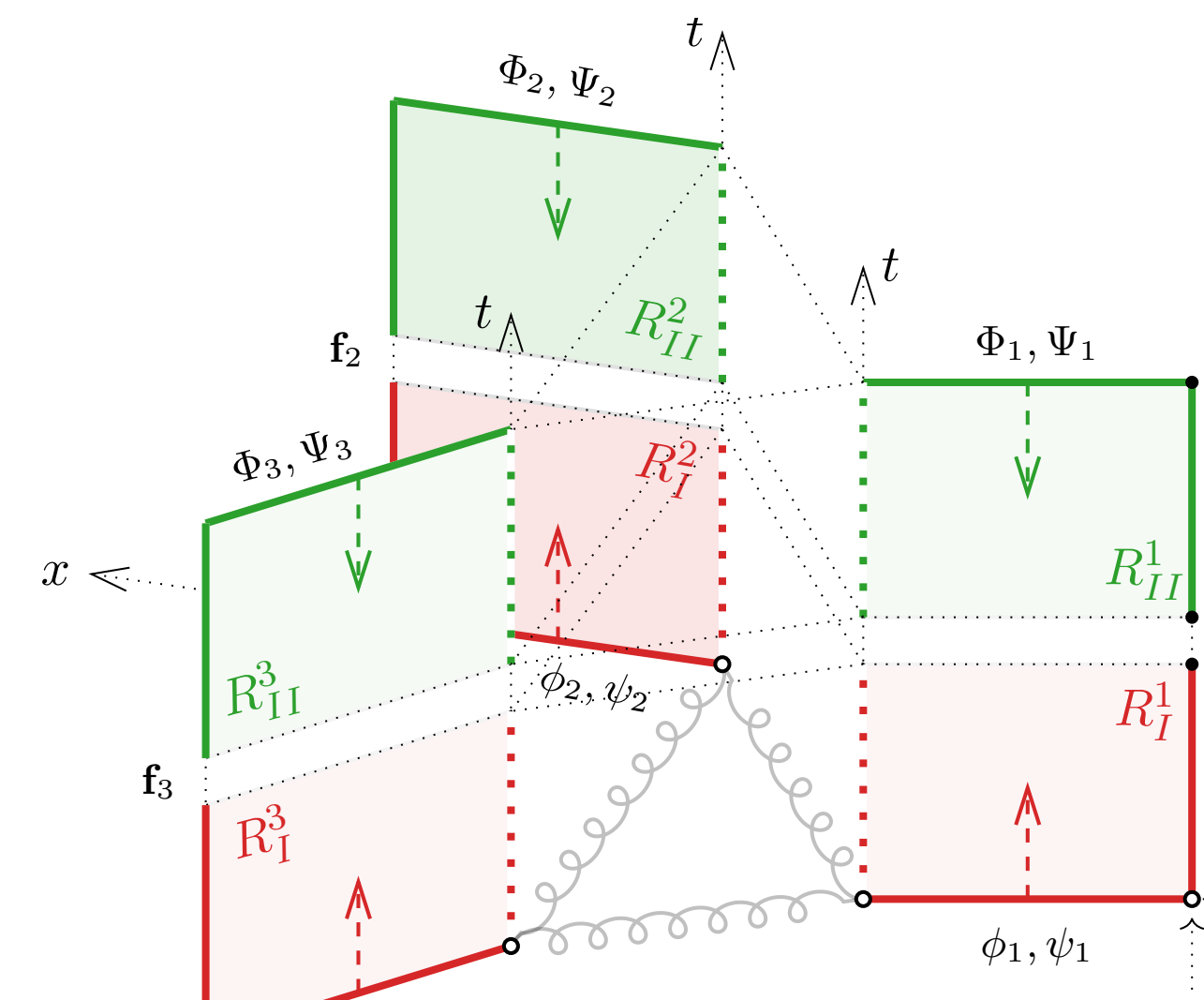
Exact boundary controllability

$$(\mathbf{E}) \begin{cases} y_{tt}^i - (K^i(y^i, y_x^i))_x = F(\mathbf{y}, \mathbf{y}_x, \mathbf{y}_t), & x \in [0, L_i], t \in [0, T] \\ y_{tt}^i(t, 0) = G^i(t, \mathbf{y}(t, 0), \mathbf{y}_x(t, 0), \mathbf{y}_t(t, 0)) \\ \quad + \int_0^t H^i(t, s, \mathbf{y}(s, 0)) ds, & t \in [0, T] \\ y^i(t, L_i) = u^i(t), & t \in [0, T] \\ (y^i, y_t^i)(0, x) = (\phi^i(x), \psi^i(x)), & x \in [0, L_i]. \end{cases}$$

The system (E) is locally exact controllable

- ▶ with n controls [G.Leugering, T.Li, Y.Wang, '18,'19].

Controllability Time (sharp):
$$T^* = \max_{i=1, \dots, n} \frac{2L_i}{\sqrt{K_{y_x}^i(0,0)}}$$



Wellposedness

We introduce $\mathbf{w}^i = (w_1^i, w_2^i, w_3^i)^T := (y^i, y_x^i, y_t^i)^T$. Then we get

$$\frac{\partial}{\partial t} \begin{pmatrix} w_1^i \\ w_2^i \\ w_3^i \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & -K_{w_2^i}^i & 0 \end{pmatrix} \frac{\partial}{\partial x} \begin{pmatrix} w_1^i \\ w_2^i \\ w_3^i \end{pmatrix} = \begin{pmatrix} w_2^i \\ 0 \\ F^i(\mathbf{w}^i) + K_{w_1^i}^i w_2^i \end{pmatrix}$$

with $(t, x) \in [0, T] \times [0, L_i]$. This, in turn, can be rewritten in the form of a quasilinear hyperbolic system

$$\mathbf{w}_t^i + A^i(x, \mathbf{w}^i) \mathbf{w}_x^i = \tilde{F}(\mathbf{w}^i),$$

where A^i has 3 distinct real eigenvalues:

$$\lambda_i^- = -\sqrt{K_{w_2^i}^i(w_1^i, w_2^i)}, \quad \lambda_i^0 = 0, \quad \lambda_i^+ = \sqrt{K_{w_2^i}^i(w_1^i, w_2^i)}.$$

Wellposedness ctd.

We may integrate the boundary conditions w.r.t. time and obtain a kind of **non-local (of time)** boundary condition in the first order system (FOS):

$$(\text{FOS}) \left\{ \begin{array}{l} \mathbf{w}_t^i + A^i(x, \mathbf{w}^i) \mathbf{w}_x^i = \tilde{F}^i(\mathbf{w}^i), \quad x \in [0, L_i], t \in [0, T] \\ w_2^i(t, 0) = \psi^i(0) + \int_0^t G^i(\tau, \mathbf{w}^i(\tau, 0)) d\tau \\ \quad + \int_0^t \int_0^\tau H^i(\tau, s, w_1^i(s, 0)) ds d\tau, \quad t \in [0, T] \\ w_1^i(t, L_i) = u^i(t), \quad t \in [0, T] \\ \mathbf{w}^i(0, x) = \mathbf{w}^{0,i}(x) = (\phi^i(x), \psi^i(x), \phi^{i'}(x)), \quad x \in [0, L_i]. \end{array} \right.$$

- ▶ Local existence of C^1 solution to (FOS) (T.Li, '85): $\exists!$ C^1 solution on $\mathcal{R}(\delta) = \{(t, x) | 0 \leq t \leq \delta, 0 \leq x \leq L\}$, where δ depends on the initial and boundary data.
- ▶ For given $T > 0$, **NO results on existence of semi-global classical solutions before.**

Lemma: A uniform priori estimate of solution to (FOS) [Y.W.'19]

$$\|w(t, \cdot)\|_1 \triangleq \|w(t, \cdot)\| + \left\| \frac{\partial w}{\partial x}(t, \cdot) \right\| \leq C(T), \quad 0 \leq t \leq T,$$

where $\|\cdot\|$ denotes C^0 -norm.

Exact boundary controllability

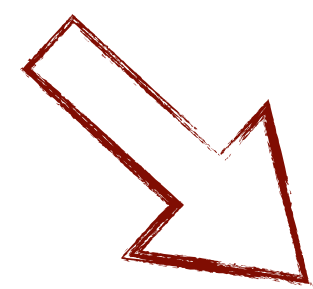
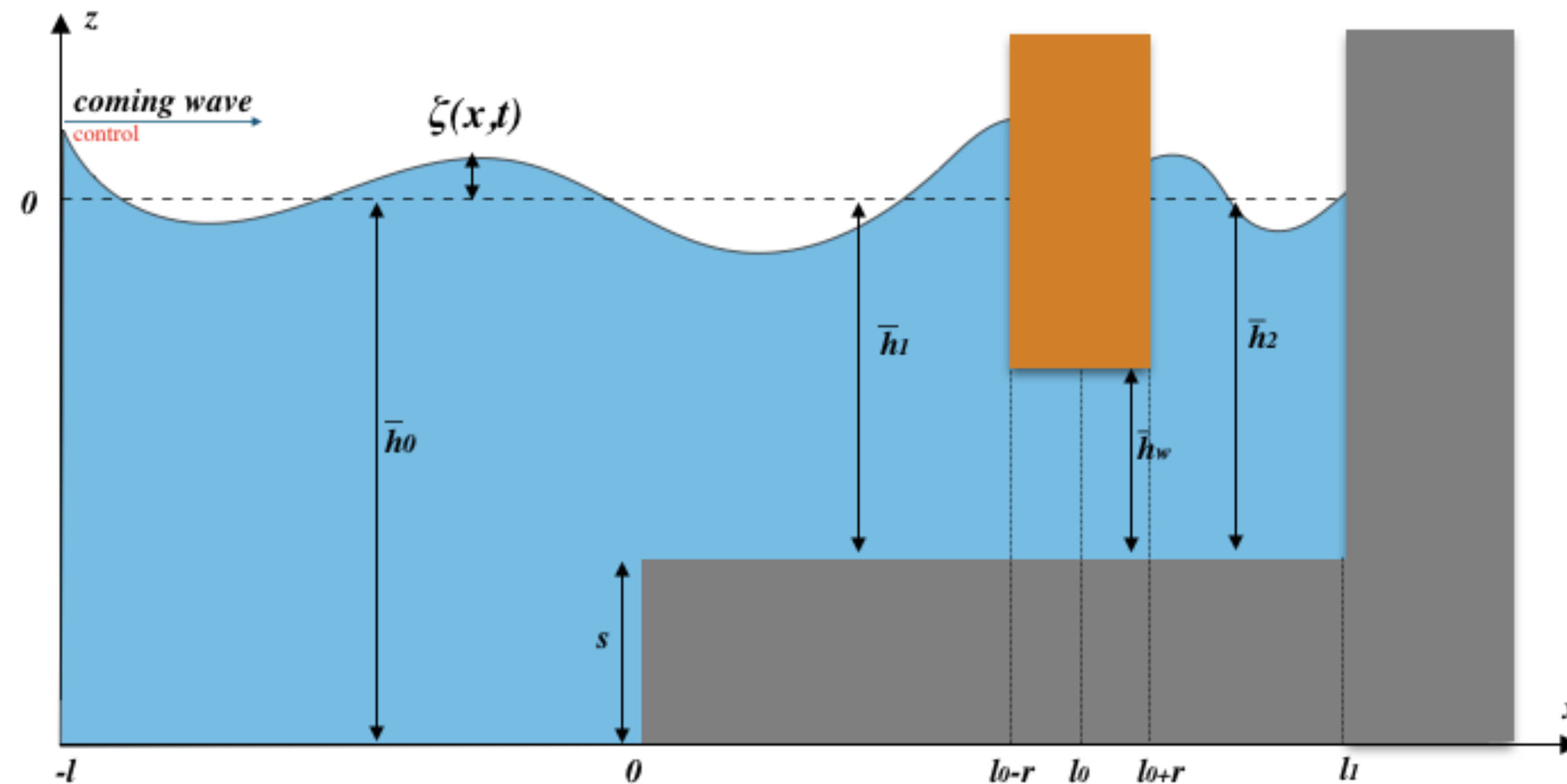
$$(\mathbf{E}) \left\{ \begin{array}{l}
 y_{tt}^i - (K^i(y^i, y_x^i))_x = F(\mathbf{y}, \mathbf{y}_x, \mathbf{y}_t), \quad x \in [0, L_i], t \in [0, T] \\
 y_{tt}^i(t, 0) = G^i(t, \mathbf{y}(t, 0), \mathbf{y}_x(t, 0), \mathbf{y}_t(t, 0)) \\
 \quad + \int_0^t H^i(t, s, \mathbf{y}(s, 0)) ds, \quad t \in [0, T] \\
 y^i(t, L_i) = u^i(t), \quad t \in [0, T] \\
 (y^i, y_t^i)(0, x) = (\phi^i(x), \psi^i(x)), \quad x \in [0, L_i].
 \end{array} \right.$$

The system (E) is locally exact controllable

- ▶ with n controls [G.Leugering, T.Li, Y.Wang, '18,'19].
- ▶ This result can be improved by **reducing the number of controls to $n - 1$** , but the space of controlled initial data is **asymmetric** [G.Leugering, S.Micu, I.Robenta, Y.Wang, '22] [G.Leugering, C.Rodriguez, Y.Wang, '22].

Extension: dynamical boundary conditions in 1st-order hyperbolic systems

- ▶ Project Conflex. [G.Vergara-Hermosilla, G.Leugering, Y.Wang, COCV '21].
- ▶ One dimensional nonlinear shallow water system, describing the free surface flow of water as well as the flow under a fixed gate structure.

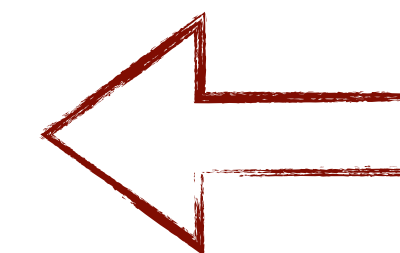


$$\begin{cases} \partial_t \zeta + \partial_x q = 0, \\ \partial_t q + \partial_x \left(\frac{q^2}{h} \right) + gh \partial_x \zeta = 0, \end{cases}$$

where $\zeta(t, x)$ is the free surface elevation, $h(t, x)$ is the fluid height, $q(t, x)$ is the horizontal discharge.

$$\zeta_0(t, 0) = \zeta_1(t, 0), \quad q_0(t, 0) = q_1(t, 0),$$

$$\begin{cases} q_2(t, l_0 + r) = q_1(t, l_0 - r) = q_w(t), \\ \left[\frac{q_i^2}{2h_i^2} + g\zeta_i \right]_{i=1, x=l_0-r}^{i=2, x=l_0+r} = -\alpha \frac{d}{dt} q_w(t), \end{cases}$$



Remark: Controllability of Nodal Profile

Nodal Control Problem: Let $T > T^* > 0$. For given **desired profile function** $y_d(t)$ to find boundary controls u^1, \dots, u^m so that

$$\mathcal{F}(u^1, \dots, u^m) = y^n(t, L_n) = y_d(t), t \in [T^*, T]$$

or $\mathcal{F}(u^1, \dots, u^m) = y_x^n(t, L_n) = y_d(t), t \in [T^*, T].$

Theorem

In a neighbourhood of an equilibrium (around 0), the system (E) is locally exact boundary controllable of nodal profile by **only 1 control** when (controllability time, sharp)

$$T > \bar{T}.$$

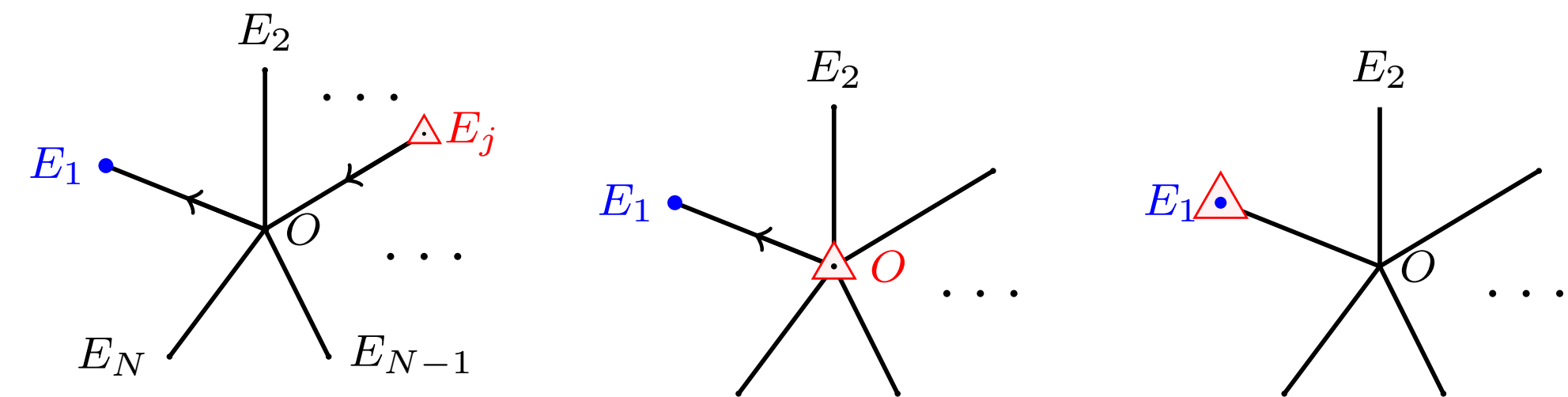
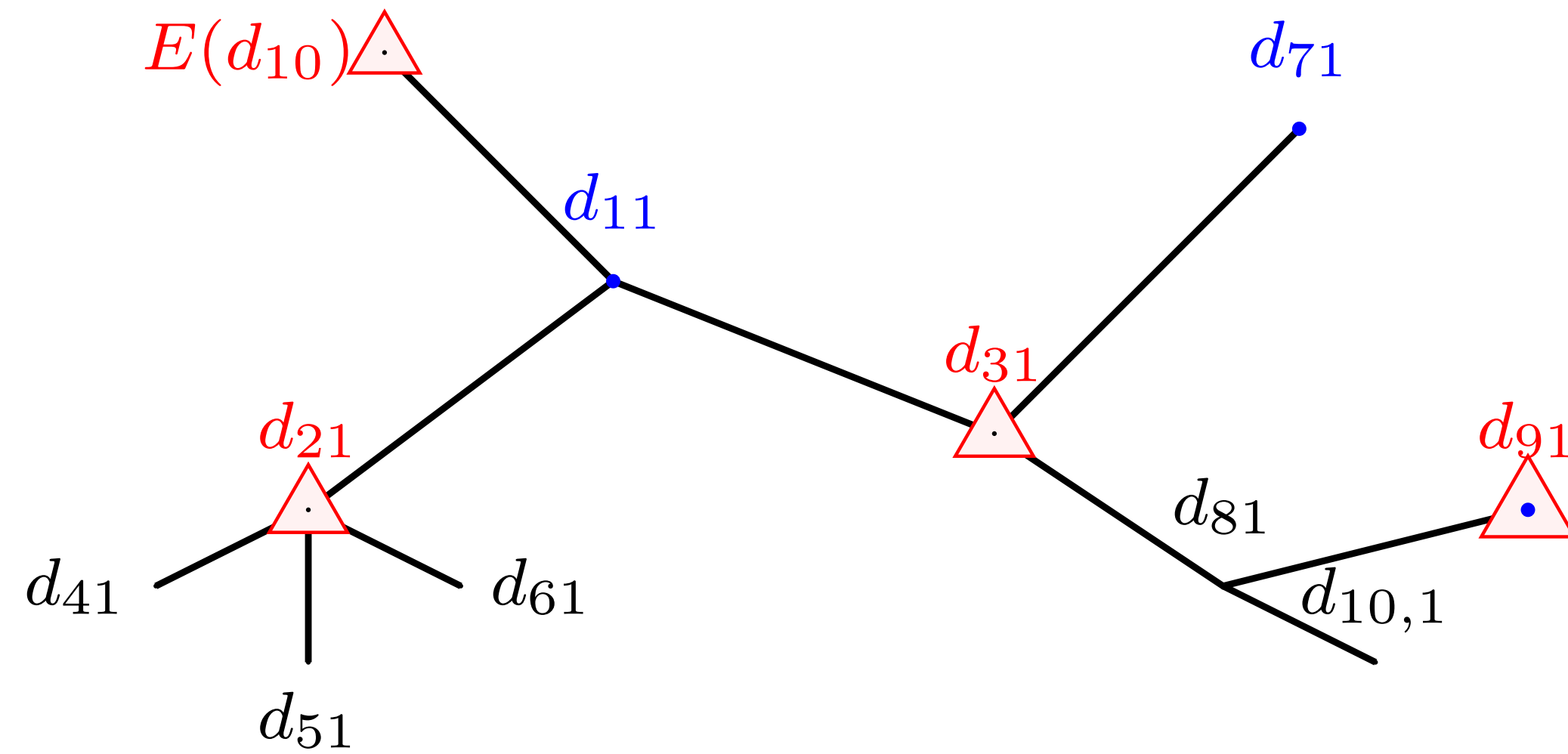


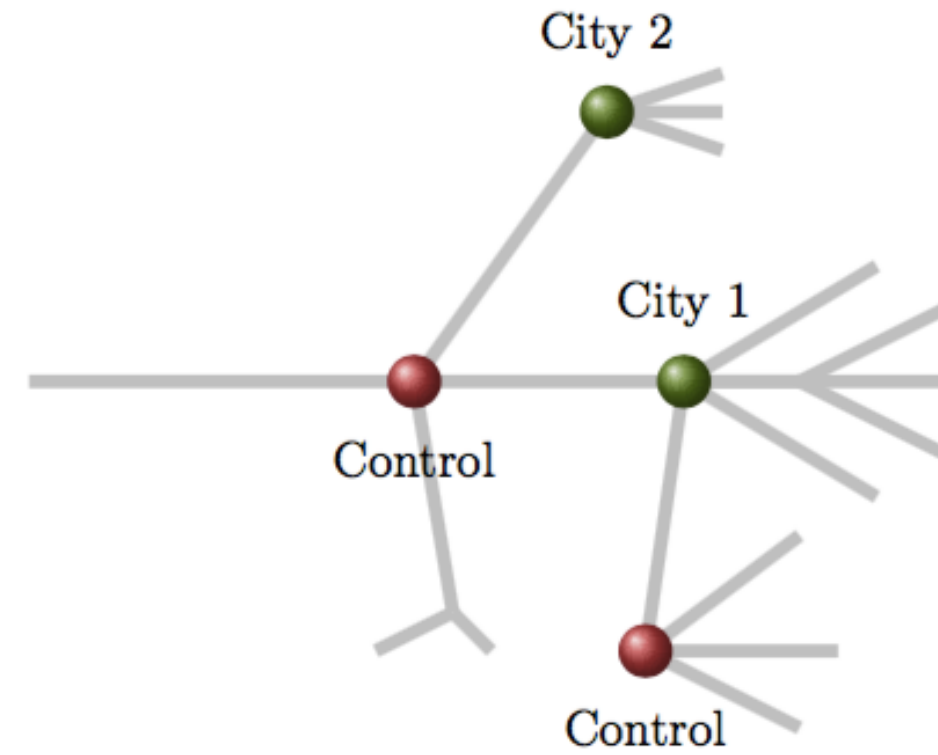
Fig.	Charged node	Controlled node	Controllability Time T
(a)	E_1	$E_j (j \neq 1)$	$\bar{T} > \frac{L_1}{\sqrt{K_{y_x^1}^1(0,0)}} + \frac{L_j}{\sqrt{K_{y_x^j}^j(0,0)}}$
(b)	E_1	O	$\bar{T} > \frac{L_1}{\sqrt{K_{y_x^1}^1(0,0)}}$
(c)	E_1	E_1 (in-situ)	$\bar{T} > 0$



- * Optimal controllability time T^* .
- * Minimum number of controls.
 - * Placement of controls.
 - * Calculation of controls.

- ▶ Nodal Profile Control: Our aim is to fit (a part of) the boundary traces to a given profile after a suitably long time $t = T$ by means of boundary controls. **[Project: Control theory on planar or spacial string networks: controllability and partial nodal control for quasilinear hyperbolic systems. (Individual funding & NSFC-1121101. Joint work with T.Li.)]**

Extension: Flow Control on gas network and Control Design.

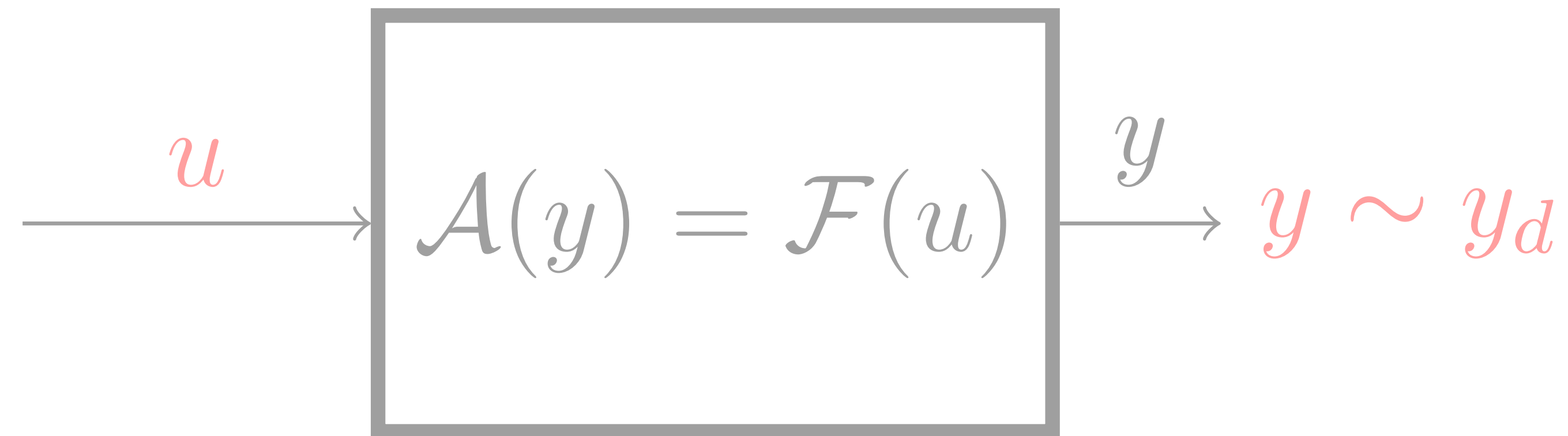


- ▶ The coupling of gas pipes.
- ▶ State function [isothermal Euler equations]
 - > $\rho(t, x)$: the density of the gas,
 - > $q(t, x)$: the flux in the pipe.
- ▶ Nodal Controls $u(t)$: Pressure increases at the compressor stations.

- ▶ **Q:** Can we find controls to satisfy the given demand of the cities?
Aim: The boundary traces of state to exactly fit any given profile as function of time on a node after a suitable time $t = T$ by means of boundary controls. [= Exact boundary controllability of nodal profile]
- ▶ **Answer:** Yes! (in local sense, and at least after a waiting time T^*).
 [M.Gugat 2010, 2014, T.Li 2010]

- * Minimum number of controls.
 - * Placement of controls.
 - * Calculation of controls.

Significant Interests

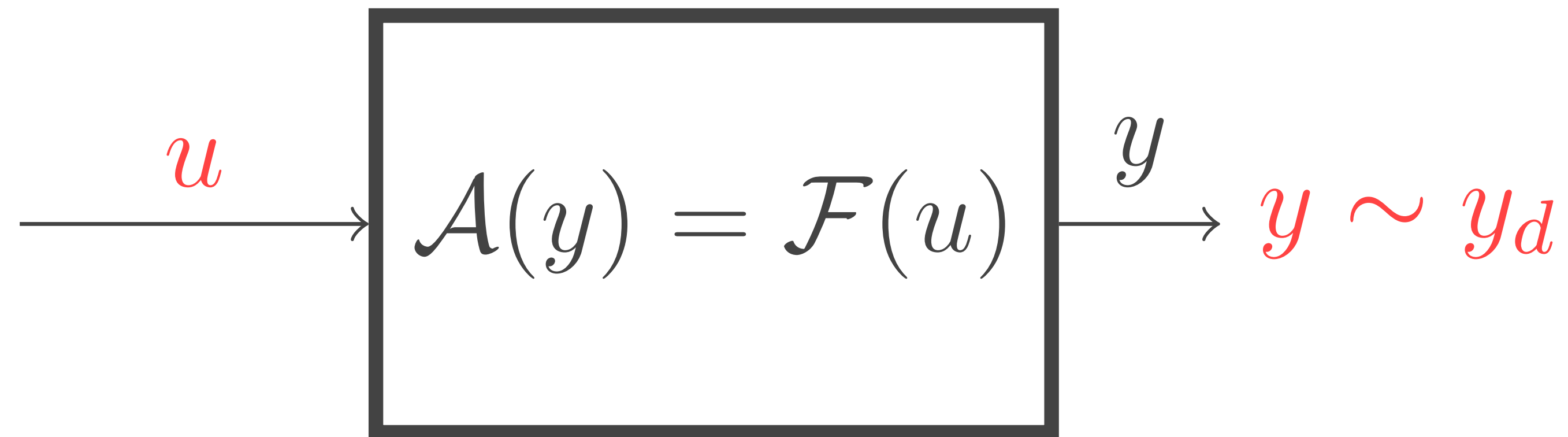


- ▶ Modeling and Analysis
- ▶ Control Theory and Optimal Design



Key issue: developing and applying mathematical methods, **including nonlinear functional analysis, new control theory and strategy** to **model, understand and control** the dynamics of PDEs.

Significant Interests

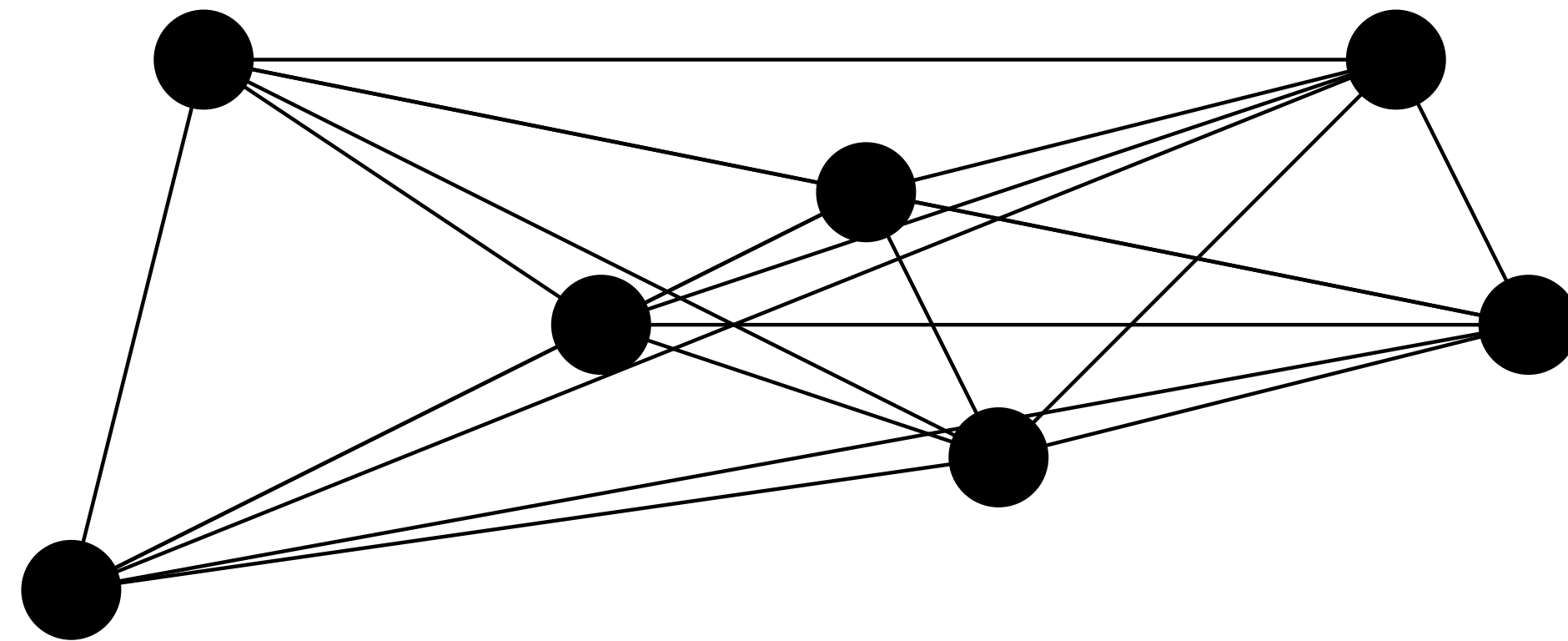


- ▶ Modeling and Analysis
- ▶ Control Theory and Optimal Design
- ▶ **Accurate and Fast Prediction of Numerical Solutions/ Optimal Control for Network PDEs**
 - 1 A stochastic method inspired by Random Batch Method
 - 2 PINN approach

A Stochastic Algorithm for the Efficient Simulation for Networked Linear hyperbolic systems

Origins of the Random Batch Method

Initial Motivation: Simulation and control of large *interacting particle systems* can be computationally demanding.

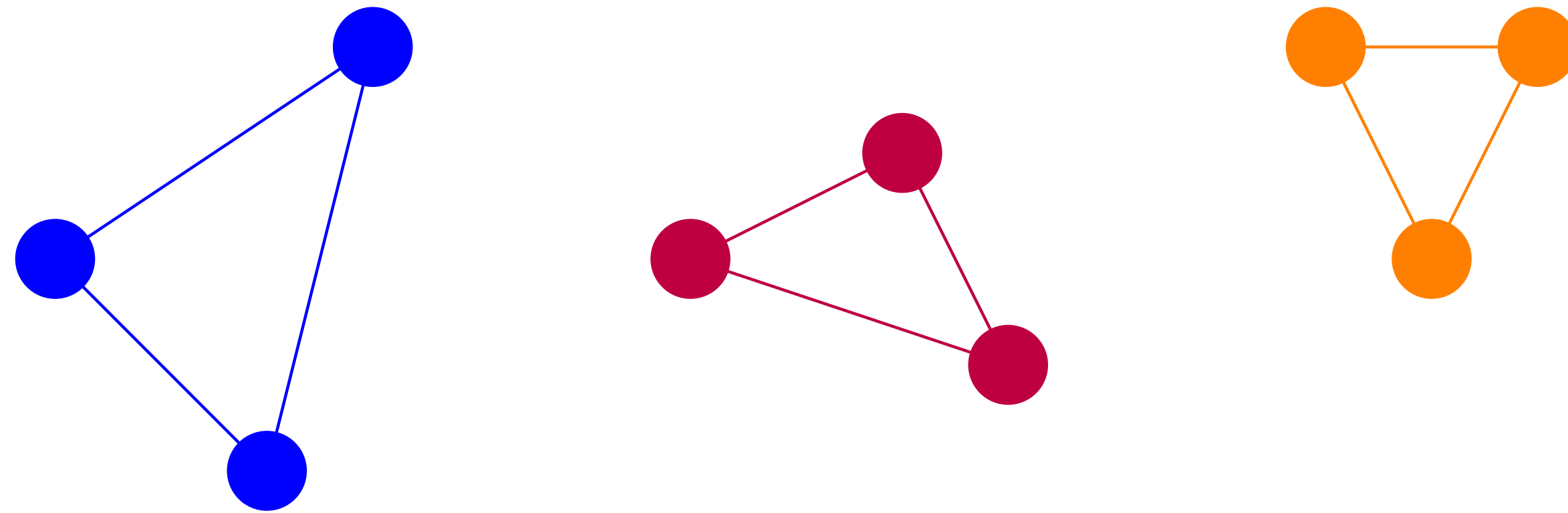


There are $N(N - 1)/2$ interaction forces between N particles.
 \Rightarrow Computational cost grows rapidly when N is large.

Origins of the Random Batch Method

Proposed simulation method: The Random Batch Method

[Shi Jin, Lei Li, Jian-Guo Liu, J. of Computational Physics, 2020]

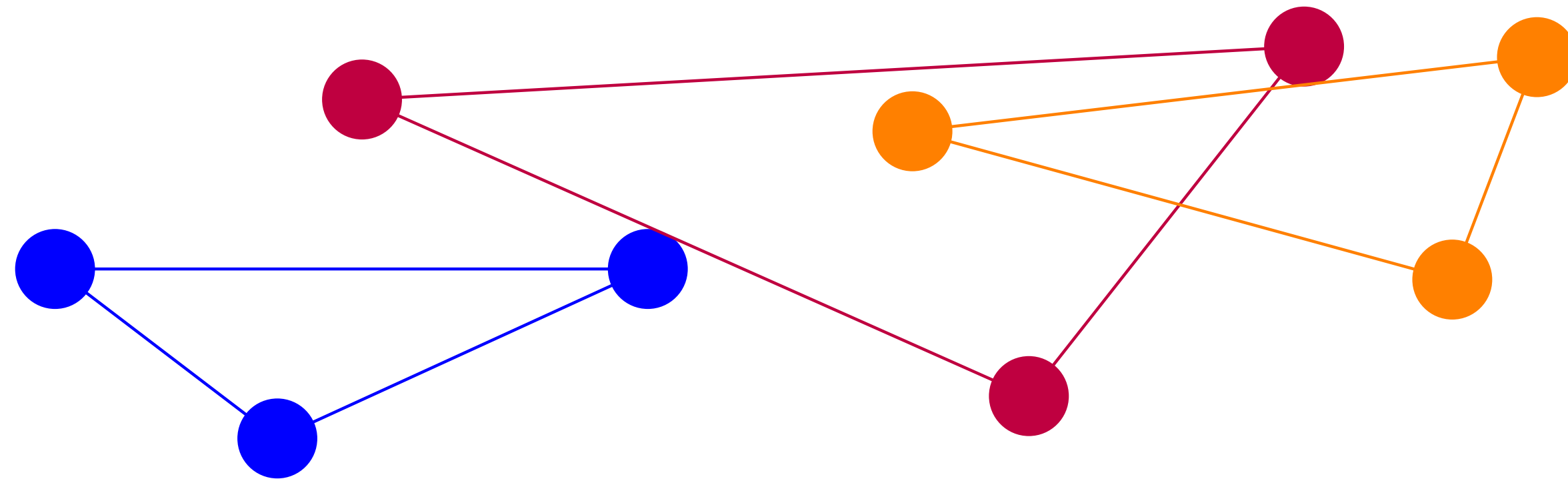


- ▶ Divide the N particles **randomly** into batches of size $P \geq 2$.
- ▶ Consider only interactions between particles in the same batch.
- ▶ Do a simulation over a short time interval of length h .

Origins of the Random Batch Method

Proposed simulation method: The Random Batch Method

[Shi Jin, Lei Li, Jian-Guo Liu, J. of Computational Physics, 2020]



- ▶ Divide the N particles **randomly** into batches of size $P \geq 2$.
- ▶ Consider only interactions between particles in the same batch.
- ▶ Do a simulation over a short time interval of length h .
- ▶ **Repeat.**

Origins of the Random Batch Method

Proposed simulation method: The Random Batch Method

[Shi Jin, Lei Li, Jian-Guo Liu, J. of Computational Physics, 2020]

- ▶ the RBM-solution converges to the solution of the original problem as $h \rightarrow 0$.
- ▶ the RBM reduces the computational cost from $O(N^2)$ to $O(PN)$.

RBM for Optimal Control

- ▶ The RBM can speed up the solution of optimal control problems governed by interacting particles systems [D. Ko, E. Zuazua, Math. Models Methods Appl. Sci., Vol. 31, No. 8, 2021] **(only numerical experiments)**.
- ▶ The first convergence proof is given in [D.Veldman, E.Zuazua, Numerische Mathematik, 2022] for finite dimensional linear-quadratic optimal control in the **operator-splitting setting**.

$$\min_u \int_0^T (|x(t) - x_d(t)|^2) + |u(t)|^2 dt,$$
$$\dot{x}(t) = Ax(t) + Bu(t), \quad x(0) = x_0$$

- ▶ Whether this algorithm can accelerate the simulation and optimization of nonlinear dynamics and **networked** infinite-dimensional systems, its convergence theory and applications are still open issues! (**A is unbounded operator**).

RBM for Hyperbolic equations: Toy Example

Consider the transport equation

$$\begin{aligned} y_t(t, x) + v(x)y_x(t, x) &= 0, & t \in (0, T), x \in \mathbb{R}, \\ y(0, x) &= y_0(x), & x \in \mathbb{R}, \end{aligned}$$

where $v(x)$ is bounded and Lipschitz, y_0 is globally Lipschitz.

- We split the generator of the semi-group as

$$\underbrace{-v(x)\frac{\partial}{\partial x}}_A = \sum_{m=1}^M \underbrace{-v_m(x)\frac{\partial}{\partial x}}_{A_m},$$

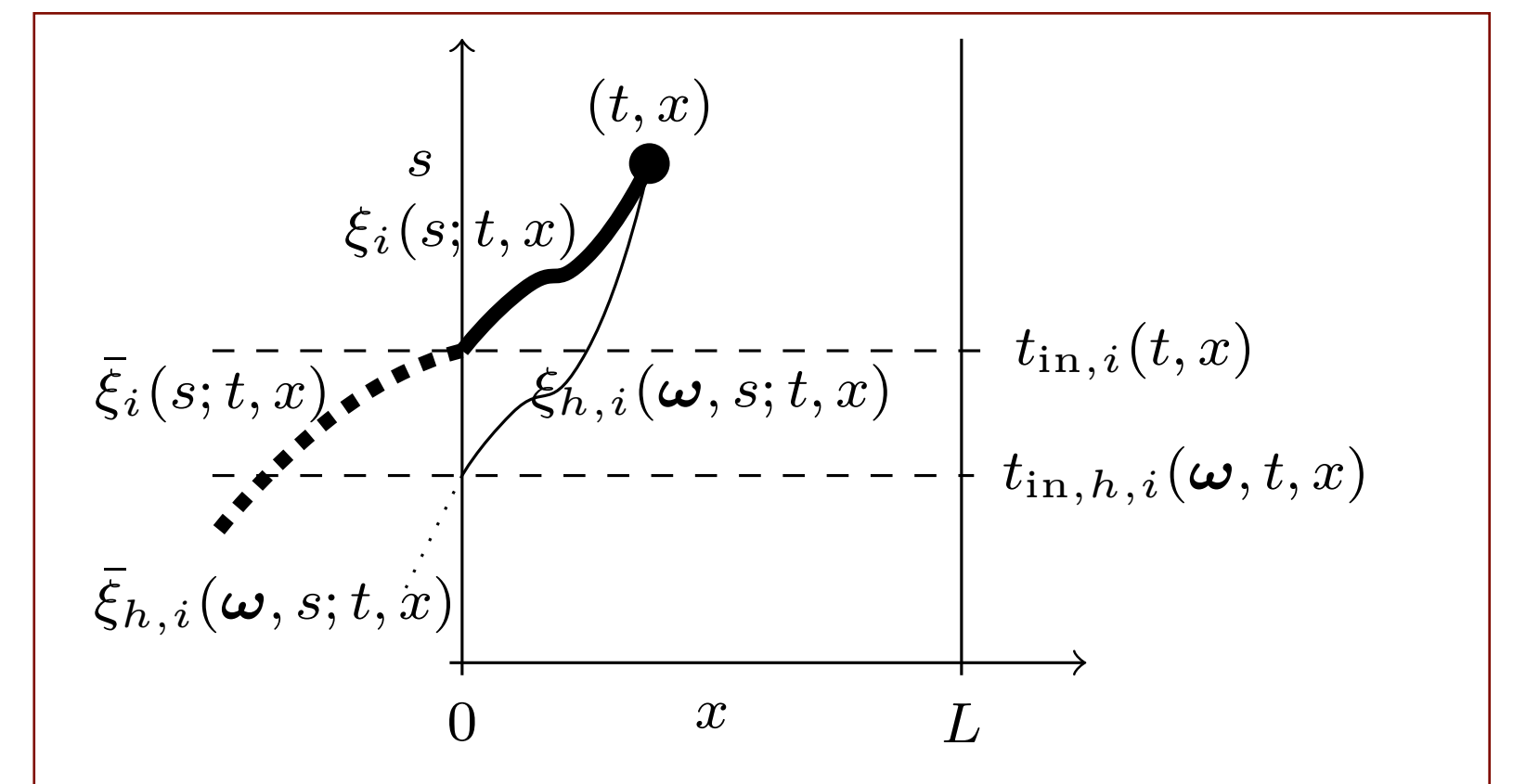
where the $v_m(x)$ are Lipschitz and bounded.

- In each time step, we randomly choose batch B_k , subset of $\{1, \dots, M\}$, of size P and consider the velocity field as

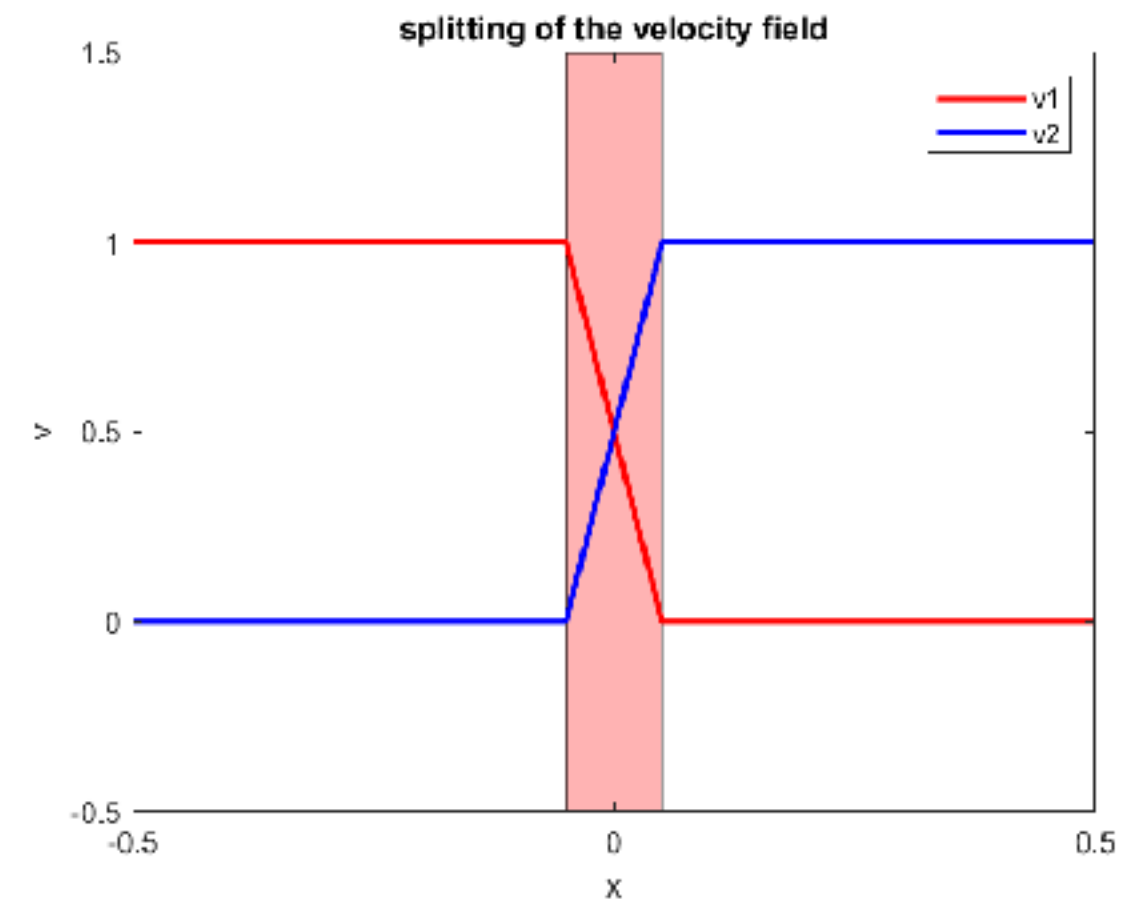
$$v_h(\boldsymbol{\omega}, x) = \frac{M}{P} \sum_{m \in B_k} v_m(x), \quad t \in [t_{k-1}, t_k).$$

Let $y_{h,t}(\boldsymbol{\omega}, t, x)$ be the solution resulting from $v_h(\boldsymbol{\omega}, t, x)$, then

$$\mathbb{E}[|y_h(t) - y(t)|_{L^\infty}^2] \leq Ch$$

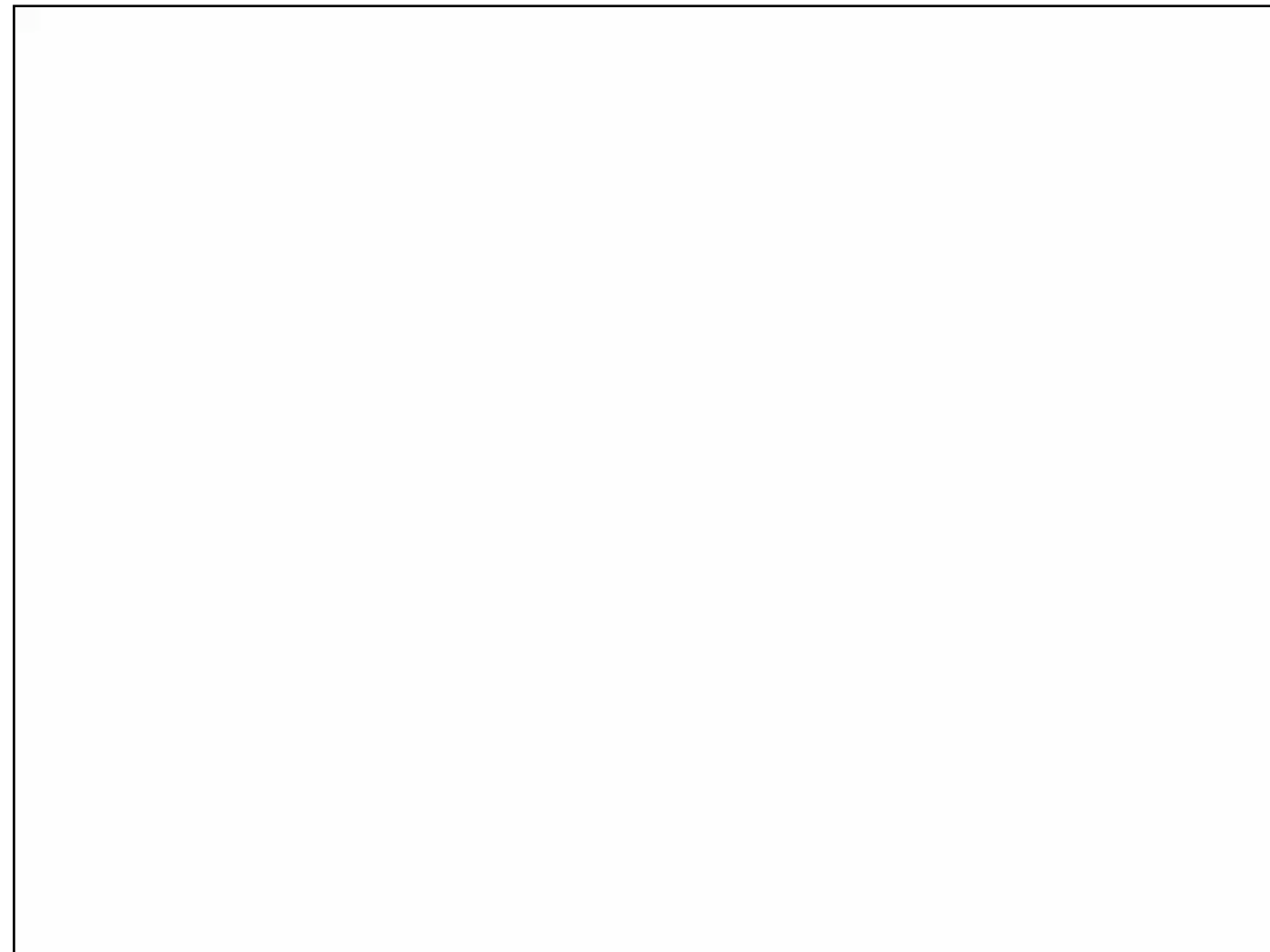


Toy example: visualization

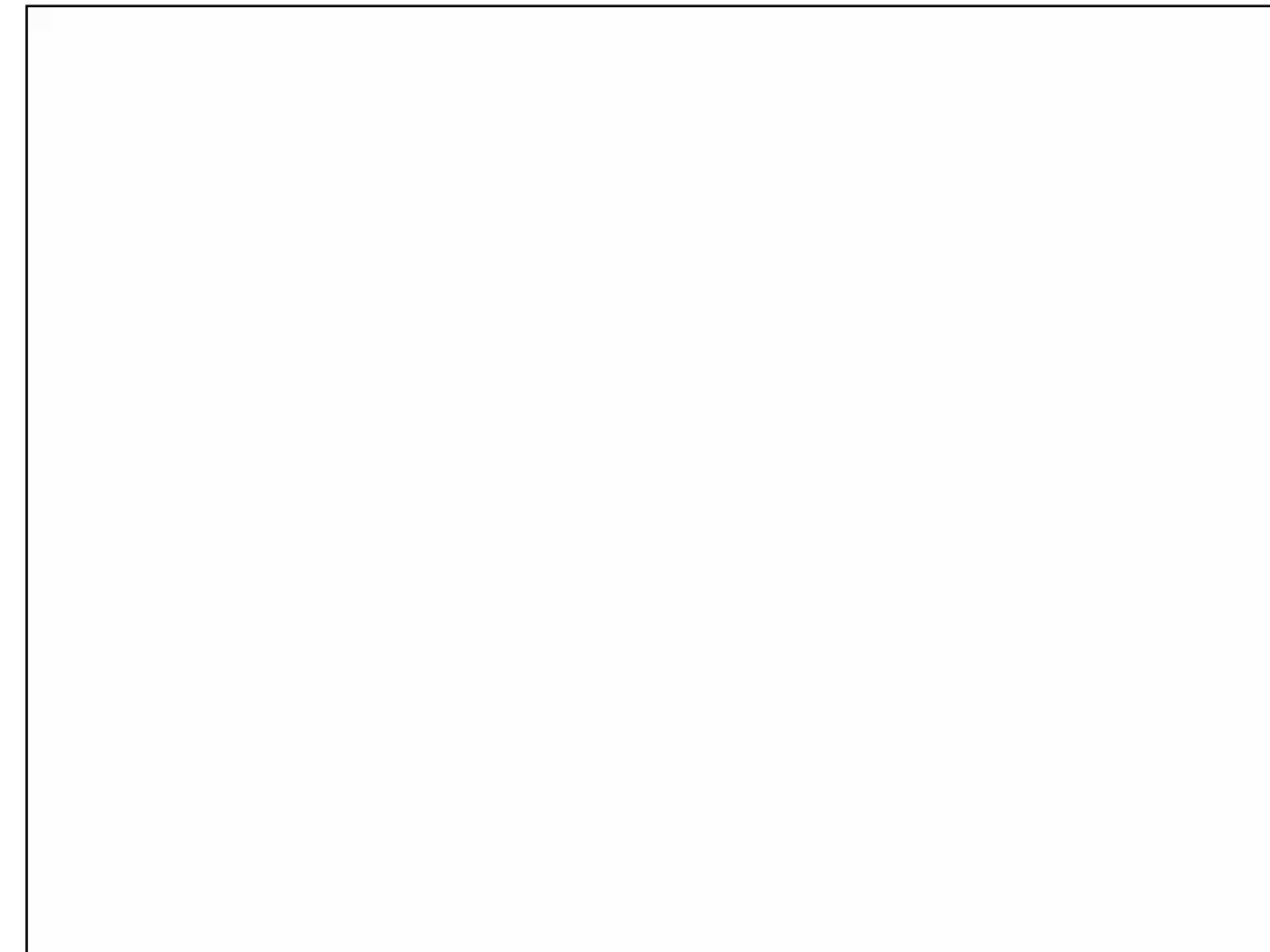


$$v(x) \equiv 1 = v_1(x) + v_2(x)$$

h=0.01

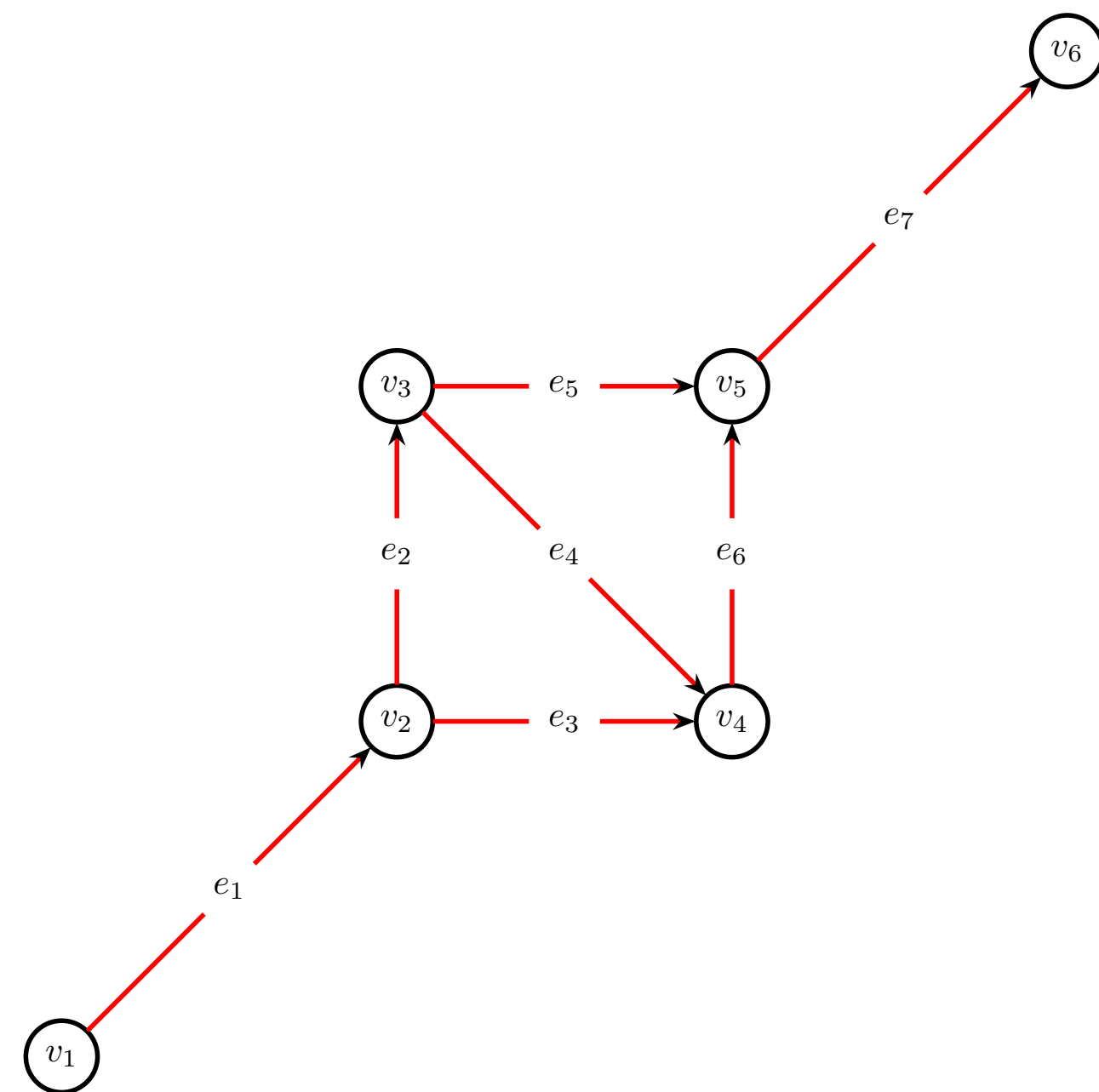


h=0.001



RBM for Networks

Coupled wave equations on diamond networks [D. W. M. Veldman, Y. Wang, 2024]



Diamond Directed Graph

$$\begin{cases} y_{tt}^{e_i}(t, x) - c_{e_i}^2 y_{xx}^{e_i}(t, x) = 0 & e_i \in E, \\ \sum_{e_i \in E(v_j)} D_{ji} c_{e_i} y_x^{e_i}(t, v_j) = \bar{u}^{v_j}(t) & v_j \in V, \\ y^{e_i}(t, v_j) = y^{e_k}(t, v_j), & \forall e_i, e_k \in E(v_j), v_j \in V, \\ y^{e_i}(0, x) = y_0^{e_i}(x), \quad y_t^{e_i}(0, x) = y_1^{e_i}(x), & e_i \in E, \end{cases}$$

$$\mathbf{w}^{e_i}(t, x) = \begin{pmatrix} w_-^{e_i}(t, x) \\ w_+^{e_i}(t, x) \end{pmatrix} = \begin{pmatrix} y_t^{e_i}(t, x) + c_{e_i} y_x^{e_i}(t, x) \\ y_t^{e_i}(t, x) - c_{e_i} y_x^{e_i}(t, x) \end{pmatrix}$$

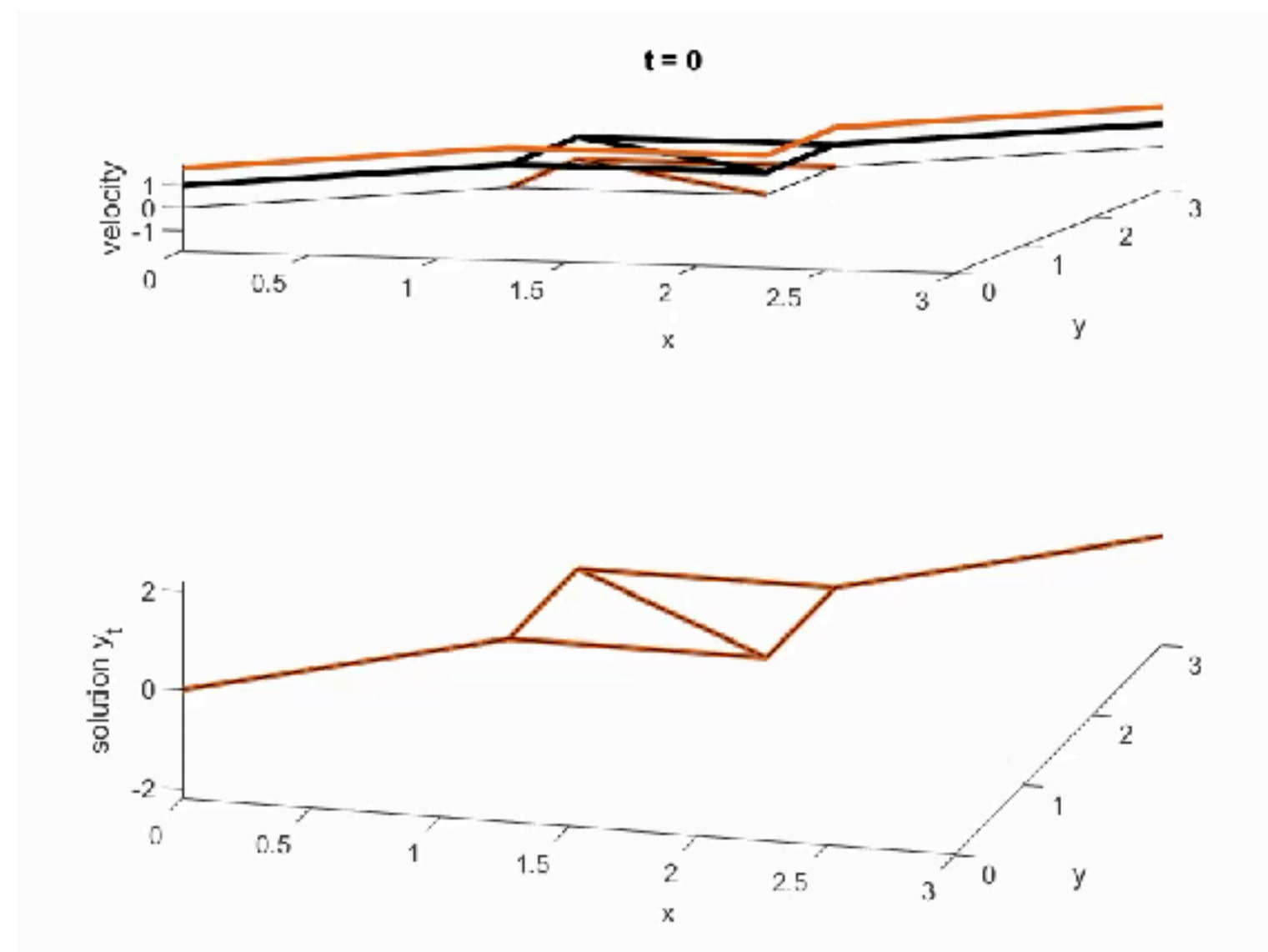
$$\begin{cases} w_{-,t}^{e_i}(t, x) - c_{e_i} w_{-,x}^{e_i}(t, x) = 0, \\ w_{+,t}^{e_i}(t, x) + c_{e_i} w_{+,x}^{e_i}(t, x) = 0. \end{cases}$$

Riemann Variables for the Wave Equation

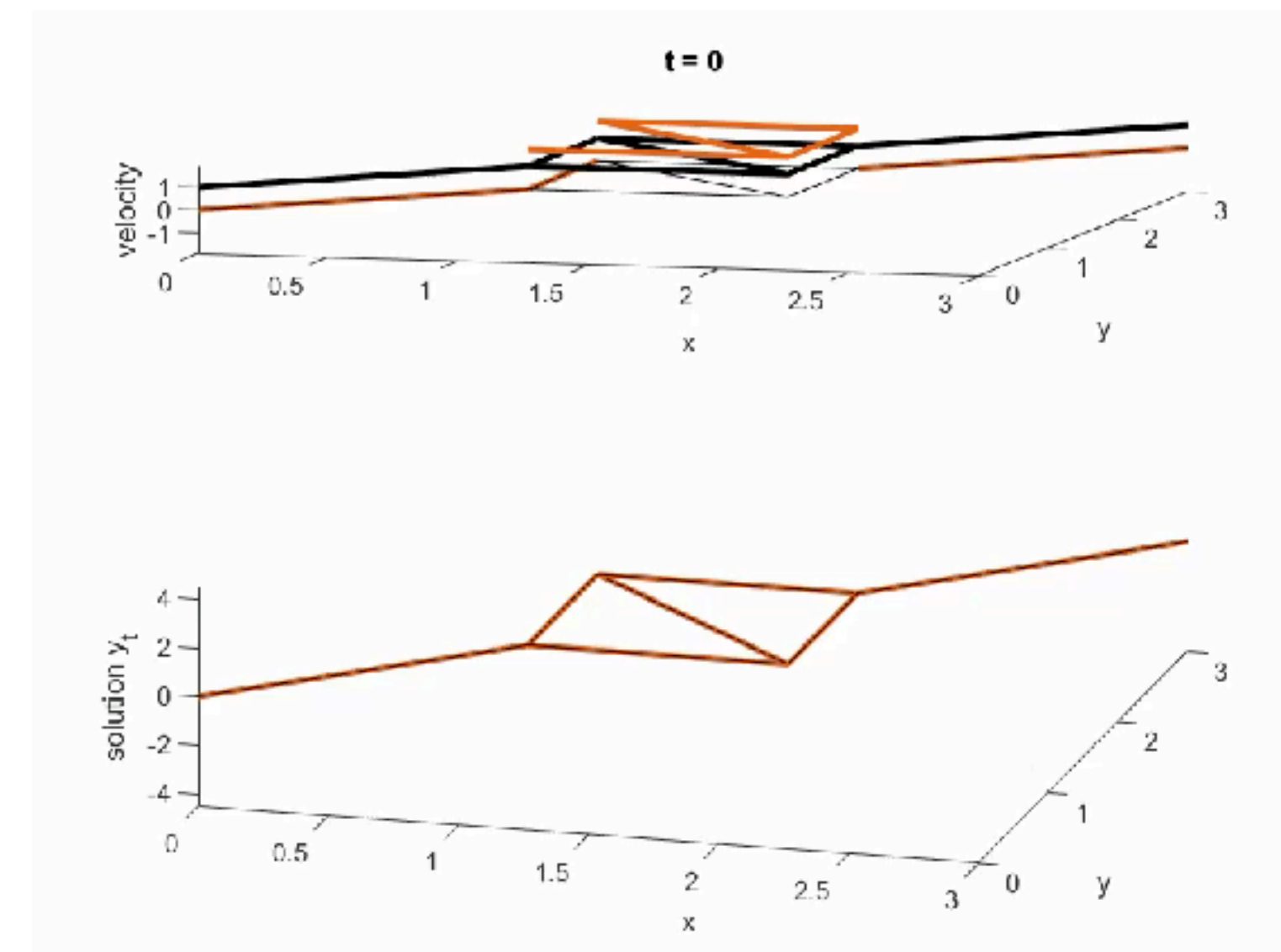
Numerical Illustration

Coupled wave equations on diamond [D. W. M. Veldman, Y. Wang, 2024]

- * Split the velocity field per edge.
- * $P = 4$ of $M = 7$ edges are active simultaneously.



- ▶ $h = 0.01, dx = 0.1$
- ▶ full Model (Black): 0.36 s
- ▶ RBM-approximation: 0.27 s
- ▶ reduction: 24%
- ▶ error: 22%



- ▶ $h = 0.001, dx = 0.01$
- ▶ full Model (Black): 551 s
- ▶ RBM-approximation: 397 s
- ▶ reduction: 28%
- ▶ error: 18%

Convergence Results

$$\min_{\mathbf{u}} J(\mathbf{u}) = \frac{1}{2} \|\mathbf{y} - \mathbf{y}_d\|_{L^2(Q)}^2 + \frac{s_0}{2} \|\mathbf{u}\|_{L^2(0,T)}^2 + \frac{s_1}{2} \|\mathbf{u}_t\|_{L^2(0,T)}^2$$

- If the original system admits an H^1 solution $y(t, x)$, then

$$\lim_{h \rightarrow 0} \mathbb{P}[|y_h^{e_i}(t) - y^{e_i}(t)|_{L^2(0, \ell_{e_i})} > \varepsilon] = 0$$

- If the original system admits an H^2 solution, then

$$\mathbb{E}[|y_h^{e_i}(t) - y^{e_i}(t)|^2] \leq Ch$$

- Remark: Markov's inequality

$$\mathbb{P}[X \geq a] \leq \frac{\mathbb{E}[X]}{a}$$

- If $s_0 > 0$ and $s_1 = 0$, then

$$\lim_{h \rightarrow 0} \mathbb{P}[|\mathbf{u}_h^* - \mathbf{u}^*|_{L^2(0,T)} > \varepsilon] = 0.$$

- If $s_1 > 0$, then

$$\mathbb{E}[|\mathbf{u}_h^* - \mathbf{u}^*|_{L^2(0,T)}] \leq Ch$$

y : original solution

y_h : solution to randomized system

\mathbf{u}^* : optimal control to original system

\mathbf{u}_h^* : optimal control to randomized system

Summary and Perspectives

- ▶ The application of the RBM to (networked) hyperbolic PDEs combines
 - > **operator splitting for PDEs**
 - > **stochastic methods for large-scale optimization**
 - > **characteristic method for 1d Hyperbolic type PDEs.**
- ▶ We **efficiently approximate** the solution to networked linear hyperbolic equations and associated optimal control problems, and obtain the **convergence results**
 - > $\mathbf{y}_h(\omega, t)$ converges to $\mathbf{y}(t)$ for $h \rightarrow 0$
 - > Convergence in the optimal controls can be proven along the lines of [E.Zuazua, D.Veldman 2022], but some regularity properties need to be verified.
- ▶ Extensions to nonlinear setting:
 - > semi-linear case is straight forward, e.g. $y_t + \Lambda y_x = f(t, x, y)$ with f Lipschitz in y .
 - > quasi-linear case is more challenging but appears in many real-world applications (nonlinear transport equations / conservation laws, and networks of incompressible Euler equations)
- ▶ Extension to non-overlapping domain decomposition on complex spatial structures and XPINNs

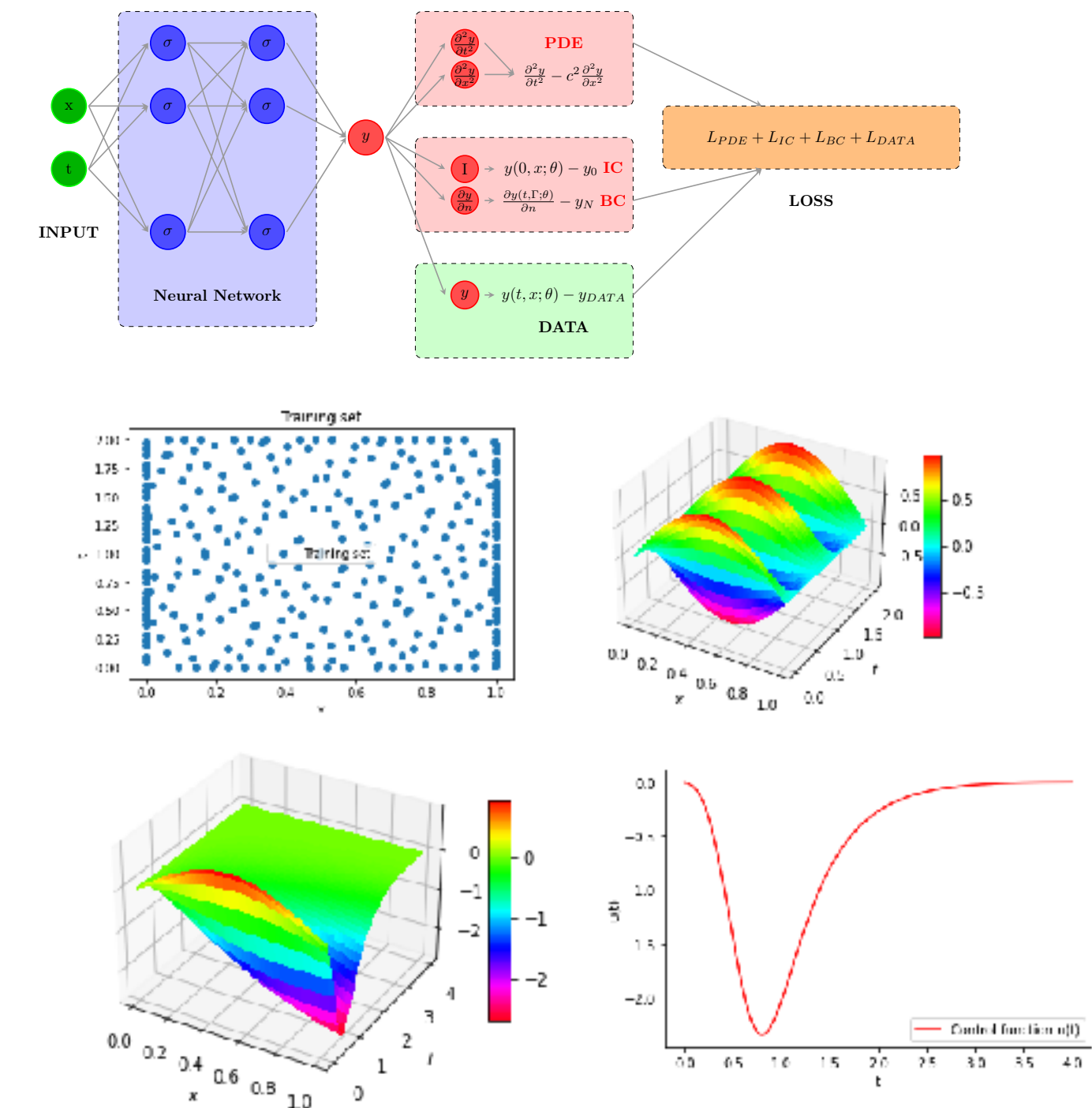
Projects on Real-time capable methods and algorithms

- **Simulation, inverse problems, and control for (degenerate) 1-D wave equations using PINNs.**

https://github.com/DCN-FAU-AvH/PINNs_wave_equation



- Dania Sana (Jun. - Sep. 2022) supervised by Y. Wang and E. Zuazua
- Internship for young female researchers at FAU-MoD (Center for Mathematics of Data)



- **PGML: Physics-Guided Machine Learning (2024-2027)** focusing on Simulation and modeling of electrochemical cells and of mechanical systems
- **SHARE at FAU (Schäffler Hub for Advanced Research at Friedrich-Alexander University)**

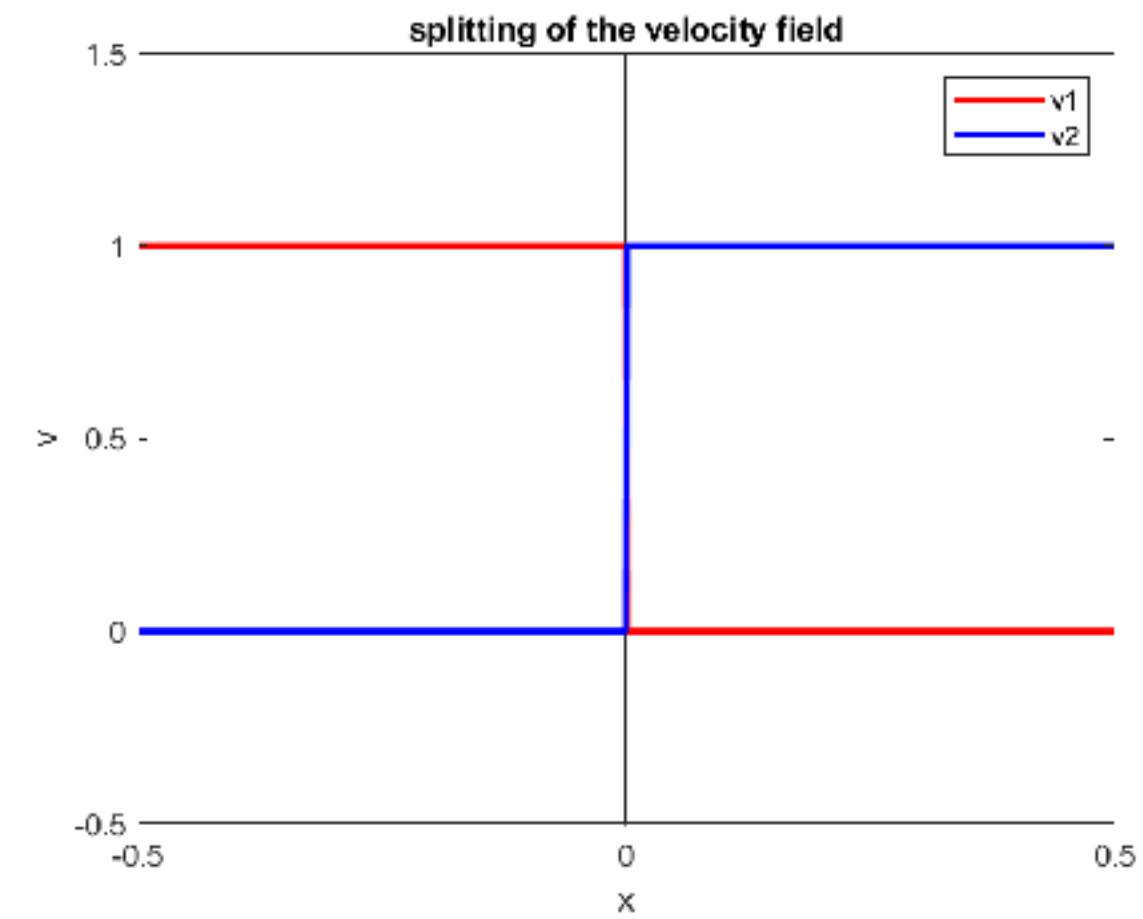


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Thank you!

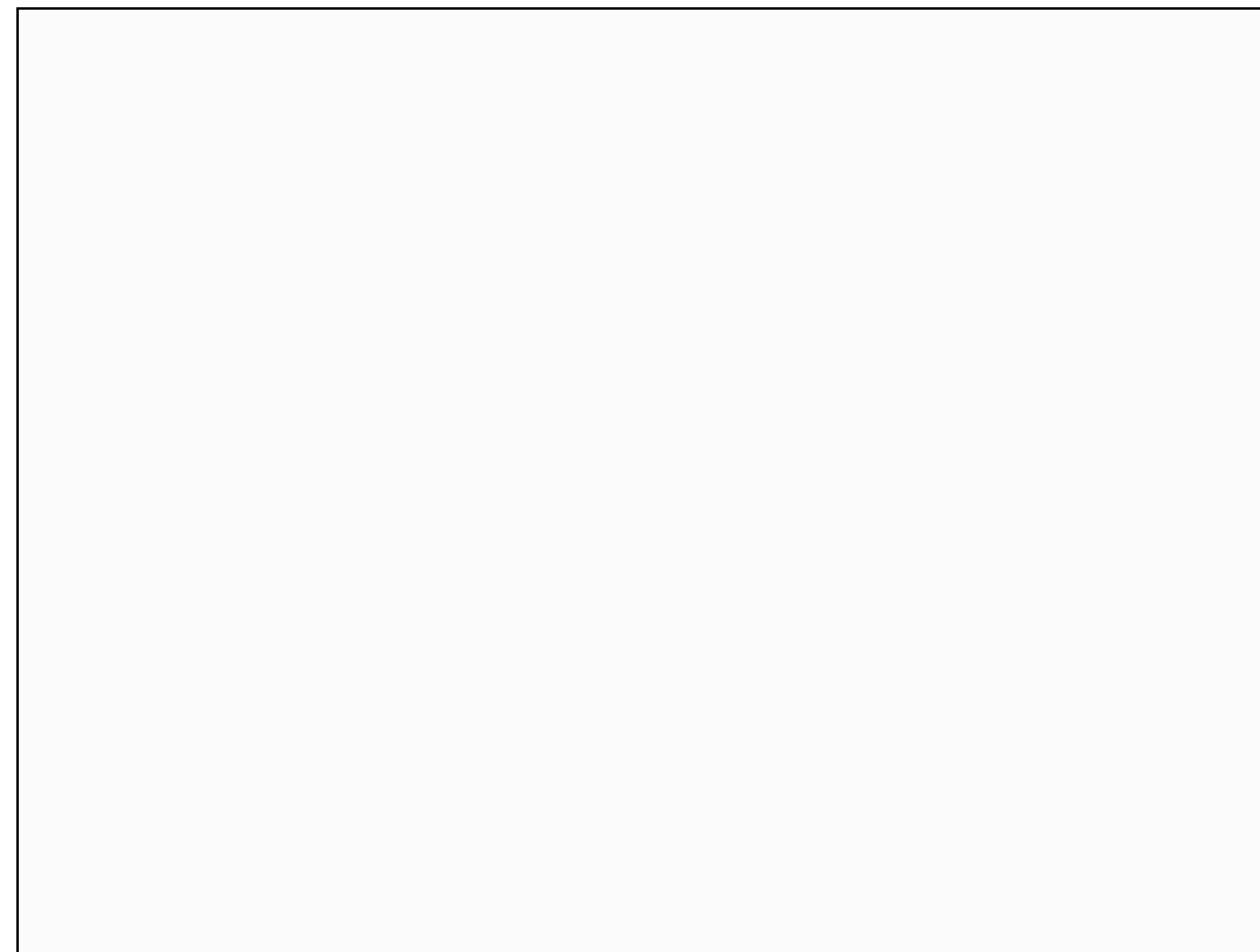
Trends in the Mathematical Sciences
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Toy example: visualization



$$v(x) \equiv 1 = v_1(x) + v_2(x)$$

h=0.001



Lemma: A uniform priori estimate of solution to (FOS) [Y.W.'19]

$$\|w(t, \cdot)\|_1 \triangleq \|w(t, \cdot)\| + \left\| \frac{\partial w}{\partial x}(t, \cdot) \right\| \leq C(T), \quad 0 \leq t \leq T,$$

where $\|\cdot\|$ denotes C^0 -norm.

Main Idea in the Proof: We apply

$$\mathbf{l}_i^- = (\mathbf{0}, \sqrt{K_{w_2}^i}, 1), \quad \mathbf{l}_i^0 = (1, \mathbf{0}, \mathbf{0}), \quad \mathbf{l}_i^+ = (\mathbf{0}, -\sqrt{K_{w_2}^i}, 1)$$

to (FOS) and define **Riemann variables** as

$$v_i = \mathbf{l}_i(w)w, \quad \bar{v}_i = \mathbf{l}_i(w)w_x.$$

They follow

$$\frac{Dv_i}{D_i t} = \sum_{j,k=1}^n \beta_{ijk}(w) \bar{v}_j \bar{v}_k + \sum_{j=1}^n \tilde{\beta}_{ij}(w) \tilde{F}_j(w) \quad (i = 1, \dots, n),$$

$$\frac{D\bar{v}_i}{D_i t} = \sum_{j,k=1}^n \gamma_{ijk}(w) \bar{v}_j \bar{v}_k + \sum_{j=1}^n \tilde{\gamma}_{ij}(w) \bar{v}_j \quad (i = 1, \dots, n),$$

along **the characteristic curves**, where

$$\frac{D}{D_i t} = \frac{\partial}{\partial t} + \lambda_i(u) \frac{\partial}{\partial x}.$$

Uniform Priori Estimate (ctd.)

Let

$$T_1 = \min_{\substack{i=1, \dots, n; \\ \|w\| \leq \eta_0}} \frac{L}{|\lambda_i(w)|} > 0.$$

For $(t, x) \in \mathcal{R}(T_1)$, we estimate $|v_i(t, x)|$ by integrating (backward) along the characteristic curve (three cases, $\lambda_i <, =, > 0$). It will arrive at $(0, \alpha)$, or (t_*, L) , or $(t_*, 0)$. In different cases, we could obtain

$$|v_i(t, x)| \leq \|v(0, \cdot)\| + C_1 \int_0^t v_i(\tau) d\tau,$$

or

$$|v_i(t, x)| \leq A \|v_i(0, \cdot)\| + \|u'\| + C_2 \int_0^t v_i(\tau) d\tau, \quad \forall t \in [0, T_1],$$

where $v(\tau) = \sup_{0 \leq t \leq \tau} \|v(t, \cdot)\|$.

Using **Gronwall inequality** it follows that

$$|v(t, x)| \leq C \max\{\|u'\|, \|v(0, \cdot)\|\} \triangleq C\alpha_0, \quad \forall t \in [0, T_1],$$

with $C > 1$.

Then **repeating** $N = \left\lceil \frac{T}{T_1} \right\rceil + 1$ times, we have

$$|v(t)| \leq C^N \alpha_0, \quad \forall t \in [(N-1)T_1, T].$$